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Observations of Magnetospheric  
Boundary Phenomena<sup>\*</sup>

by

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## ABSTRACT

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A brief summary of recent major observations of magnetospheric boundary phenomena is provided with emphasis on Explorers 12 and 14 and IMP 1 measurements of charged particles and of magnetic fields in the sunward magnetopause region and within the tail of the magnetosphere in the anti-solar direction. Major observations of charged particles in these regions are summarized in several equatorial spatial distribution diagrams.

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## I. INTRODUCTION

From early observations of auroras [Störmer, 1955] and of cometary tails [Biermann, 1951, 1957] the earth was thought to be immersed in a plasma streaming radially outward from the sun. On the basis of observations of cometary tails and of the solar corona a hydrodynamic model of the solar corona was developed [Parker, 1958] which later proved to be in remarkable agreement with direct measurements [Snyder, Neugebauer, and Rao, 1963] of the solar corpuscular radiation, or solar wind. The advent of the artificial earth satellite and the interplanetary probe provided in situ measurements of the environment of the earth in the interplanetary medium: a continual solar plasma directed radially outward from the sun with a bulk velocity of  $300\text{--}600 \text{ km (sec)}^{-1}$  (corresponding proton kinetic energy of  $\sim 1 \text{ keV}$ ), a density of  $\sim 1$  to  $20 \text{ protons (cm)}^{-3}$ , and a characteristic temperature of  $\sim 10^5\text{--}10^6 \text{ }^\circ\text{K}$  [Snyder et al., 1963; Bridge et al., 1965], threaded by a weak interplanetary field of  $\sim 5\gamma$  ( $1\gamma = 10^{-5} \text{ gauss}$ ) forming a general Archimedean spiral structure stretching outward from the sun [Ness and Wilcox, 1965]. The motion of the solar wind past the earth is impeded by the presence of the earth's magnetic field. This interaction of the solar plasma with the geomagnetic field is presumed to be directly related, for example, with the auroras, the geomagnetic storm, and the energetic charged particles of the earth's radiation zones. The solar

wind confines the earth's magnetic field to a cavity, or magnetosphere, which is terminated in the solar direction at  $\sim 10 R_E$  ( $R_E$  = earth radius) geocentric radial distance [Cahill and Amazeen, 1963; Ness et al., 1964] and is drawn into a magnetic tail extending beyond  $30 R_E$  in the antisolar direction [Ness, 1965] as suggested qualitatively by earlier theoretical models [cf. Piddington, 1959, 1960; Johnson, 1960; Axford and Hines, 1961]. Since the propagation of disturbances (Alfvén waves) in the interplanetary medium can proceed at only a fraction ( $\sim 1/5$ ) of the bulk velocity of the solar plasma, a collisionless magnetohydrodynamic shock should develop on the sunward side of the magnetosphere [cf. Axford, 1962; Kellogg, 1962] and has been directly observed [Ness et al., 1964; Freeman, 1964; Bridge et al., 1965]. Measurements of charged particles near and in the outer magnetosphere have revealed reservoirs of charged particles energized by 'local' acceleration mechanisms driven by the solar wind, for example in the transition region between the shock wave and the sunward magnetospheric boundary and also in the tail of the magnetosphere, which are indicative of charged particle sources for auroral phenomena. The energy flux supplied by the solar wind to the area of the magnetosphere projected toward the sun and available to drive magnetospheric and upper atmospheric phenomena is  $\sim 10^{20}$  ergs (sec) $^{-1}$ , and is sufficient for supplying the average corpuscular energy flux  $\sim 10^{18}$  ergs (sec) $^{-1}$  precipitated into the earth's

upper atmosphere at auroral latitudes [cf. O'Brien, 1964; Frank and Van Allen, 1964a]. In the following presentation the major results of an already substantial and interesting collection of experimental data concerning phenomena in the earth's outer magnetosphere (a complex and variable structure with a volume of at least  $10^5 R_E^3$ ) and its immediate vicinity will be discussed.

## II. COORDINATE SYSTEMS

A coordinate system for the description of the geomagnetic field direction and magnitude and the spatial and angular distributions of charged particles within the magnetosphere is necessary for the understanding and concise summary of experimental data. Examples of extremely useful coordinates for the summary and analysis of charged particle distributions in the magnetosphere are the familiar B and L coordinates [McIlwain, 1961] based upon the adiabatic invariants of charged particle motion in a magnetic field and measurements of the geomagnetic field at the surface of the earth. This coordinate system provides a two-dimensional map of a three-dimensional spatial distribution of charged particles and temporal changes and longitude nondegeneracies, for example, in the B-L map can be interpreted in terms of the violation of the adiabatic invariants and of sources and sinks of charged particles. Since the flow of solar plasma perturbs the distant geomagnetic field so that an extrapolation of the surface measurements of the geomagnetic field (assuming  $\nabla \times \vec{B} = 0$  everywhere external to the earth's surface) becomes inaccurate it is to be expected that this surface-derived B and L coordinate system becomes invalid for the organization of particle intensities in the distant magnetosphere. This situation also becomes more acute as predominant magnetic control of the motion of the charged particle is subdued in the distant magneto-

sphere by electric drifts, by rapid radial diffusion, and by loss and source mechanisms effective in periods less than a longitudinal drift period. Realistic models of the distorted geomagnetic field have been introduced [cf. Hones, 1963; Mead, 1964] and the motions of charged particles in these models [Malville, 1960; Fairfield, 1964] are studied by utilizing the adiabatic invariants in order to separate, for example, the effects of gradient drifts from electric drifts [Hones, 1963; Taylor and Hones, 1965; Taylor, 1965]. Significant perturbation of the observed geomagnetic field from the field as derived from a geomagnetic potential based upon surface measurements begins at 6 to  $7 R_E$  in the magnetic equatorial plane [Cahill and Amazeen, 1963; Ness et al., 1965]. A significant perturbation of the geomagnetic field by the solar wind beyond  $L \approx 6$  has also been reflected in the measurements of energetic electrons  $E \sim 1.5$  MeV by showing that (a) a nondegeneracy of intensities at a given  $L$  as a function of local time [Frank, 1965b] and (b) an asymmetry in observed omnidirectional intensities about the magnetic equator [Frank, 1965a] exist when the surface-derived geomagnetic field is used for the determination of  $B$  and  $L$  of the positions of the observations. Since the solar wind is responsible for the formation of the magnetosphere (terrestrial ring currents, closed or open, and current sheets will also perturb the distant geomagnetic field) a coordinate system based upon the solar



direction immediately appears appropriate. Such coordinates were first used in the evaluation of Explorer 10 observations of plasma [Bonetti et al., 1963] and magnetic fields [Heppner et al., 1963] and were designated as solar ecliptic coordinates. In the top diagram of Figure 1 is an illustration of this right-handed Cartesian coordinate system ( $X_{SE}, Y_{SE}, Z_{SE}$ ):  $\vec{X}_{SE}$  is parallel to the earth-to-sun direction and  $\vec{Z}_{SE}$  is directed toward the North ecliptic pole,  $\theta_{SE}$  and  $\phi_{SE}$  are the solar ecliptic latitude and longitude, respectively, of point P. For example, in an earth-centered solar ecliptic coordinate system the solar wind vector is directed toward  $\theta_{SE} = 0^\circ$ ,  $\phi_{SE} = 180^\circ$  neglecting the aberration effect of  $\sim 5^\circ$  due to the earth's revolution about the sun. A more recent coordinate system based upon IMP 1 magnetic measurements in the tail of the magnetosphere has been developed [Ness, 1965] which recognizes the role of the direction of the earth's magnetic moment as well as the solar direction. These coordinates are referred to as solar magnetospheric coordinates ( $X_{SM}, Y_{SM}, Z_{SM}$ ) (refer to bottom, Figure 1) and are defined with  $\vec{X}_{SM}$  parallel to  $\vec{X}_{SE}$  and  $\vec{Y}_{SM}$  parallel to  $\vec{M}_E \times \vec{X}_{SM}$  where  $\vec{M}_E$  is parallel to the earth's magnetic dipole axis. In an earth-centered solar magnetospheric coordinate system, for example, a point fixed in celestial coordinates will move in an annual cyclic motion due to the earth's revolution about the sun upon which is superimposed a diurnal motion due to the rotation of the earth. Extensive

surveys of magnetic fields and particle distributions in the tail of the magnetosphere are currently the indicators as to whether magnetic, solar ecliptic, or solar magnetospheric coordinates provide the highest degree of coherency in the organization of the data. A summary of the major experimental observations directed toward this topic is as follows:

- (1) The neutral sheet in the earth's magnetic tail is more precisely defined by the solar magnetospheric equatorial plane [Ness, 1965].
- (2) Vela satellite measurements of electrons  $E > 50$  keV at  $\sim 17 R_E$  show that the occurrence of measurable intensities is preferentially near the magnetic equator rather than near the ecliptic plane or solar magnetospheric equatorial plane [Montgomery et al., 1965].
- (3) Explorer 14 measurements of the omnidirectional intensities of electrons  $E > 40$  keV near the local midnight meridian indicated a tail of electron intensities lying more nearly in the ecliptic plane than near the geomagnetic equator beyond  $10 R_E$ . In fact measurements during a relatively quiescent period in magnetic activity and in intensity temporal variations at similar geomagnetic latitudes but at different ecliptic latitudes repeatedly indicated that geomagnetic latitude was inadequate for

organizing the observed intensities [Frank, 1965a].

- (4) The further sophistication of the solar magnetospheric coordinates has not yet proved to be necessary for the organization of shock wave and magnetopause positional data. Solar ecliptic coordinates continue in practice to be adequate for the summary of these data [cf., Ness et al., 1964; Bridge et al., 1965].

It must be remembered in consideration of the above results that the observation of electron intensities  $E > 50$  keV in the tail region (for example, (2) above) is not necessarily indicative of the coincident presence of the neutral sheet (1). Further surveys of the electrons and protons in this region of the magnetosphere are necessary in order to pursue the interrelationships of the above results.

### III. MEASUREMENTS OF THE DISTANT GEOMAGNETIC FIELD

It is easily recognized that the interpretation of charged particle observations in the earth's magnetosphere and its vicinity depends heavily upon the character of the associated magnetic fields. Major experimental evidences concerning the topology of these magnetic fields are summarized here.

Magnetopause-transition-shock regions. One of the early direct observations of the interaction of the solar wind plasma with the geomagnetic field is shown in Figure 2 as Explorer 12 passed through the magnetospheric boundary, or magnetopause, near the earth-to-sun direction [Freeman, Van Allen, and Cahill, 1963]. The scalar value  $|\vec{F}|$  of the observed magnetic field (center, Figure 2) monotonically increases over the Finch and Leaton (1955) extrapolated field with increasing geocentric radial distance while remaining relatively constant in direction. At 52,000 km ( $8.2 R_E$ ) with  $|\vec{F}|$  approximately a factor of 2 larger than the theoretical field a sharp decrease in the magnitude and change of character of the magnetic field occurred and highly disordered magnetic fields with magnitudes of tens of gammas were observed at larger radial distances. Simultaneous observations of charged particles displayed (1) a severe decrease of trapped electron  $40 \text{ keV} < E < 50 \text{ keV}$  intensities (SPL-SPB, top, Figure 2) and (2) a large increase of low-energy electrons  $\sim$  a few keV of energy

with increasing radial distance concurrent with the above change in the character of the magnetic field. These measurements were interpreted as (1) the termination of durable geomagnetic trapping of charged particles at  $8.2 R_E$  (2) concurrent with the termination of a perturbed but smooth geomagnetic field at the magnetopause with (3) the peak of quasi-thermalized plasma beyond the magnetopause supplying the plasma pressure necessary to balance the magnetic pressure due to the compressed geomagnetic field just inside the magnetopause. These early observations, although not the most detailed available at present, provide a valuable guide to the gross characteristics of the sunward magnetopause region. Although the radial thickness of the persistent quasi-thermalized plasma observed with Explorer 12 was typically 2 to 3  $R_E$  in depth [Freeman, 1964] in agreement with theoretical predictions of the stand-off distance of the bow shock [Kellogg, 1962; Axford, 1962], no observations with the Explorer 12 magnetometer of a change in the character of the magnetic field designating the position of the shock wave were reported. Direct observations of a collisionless magnetohydrodynamic shock wave with a magnetometer were first reported by Ness et al. [1964] with the IMP 1 satellite. An example of these observations is given in Figure 3 which displays the magnetopause at  $13.6 R_E$  and the disordered magnetic field at greater radial distances in the region between the magnetopause and the shock wave. The coordinates  $\theta$ ,  $\varphi$ ,  $X_{SE}$ ,  $Y_{SE}$ , and  $Z_{SE}$

are satellite-centered solar ecliptic coordinates. At a radial distance of  $19.7 R_E$  and beyond the variances in the magnetic field components of discrete sets of measurements ( $\delta X_{SE}$ ,  $\delta Y_{SE}$ ,  $\delta Z_{SE}$ ) become small relative to those within the transition region (or 'magnetosheath') at smaller radial distances and a similar effect is noticeable in the averaged scalar field  $\overline{F}$ . The change in the character of the magnetic field at  $19.7 R_E$  which is a persistent feature of the IMP 1 measurements beyond the sunward side of the magnetosphere, although variable in position, is interpreted by Ness et al. [1964] as delineating the shock wave boundary. Ness et al. [1964] have summarized their crossings of the magnetopause and shock wave in the ecliptic plane by rotating the observations in a meridian plane containing the  $Z_{SE}$  axis and the position of the observation; this summary is reproduced in Figure 4. These data have also been normalized, or rectified, to a geomagnetic latitude of the subsolar point  $X_{SS} = 0^\circ$ . In the ecliptic plane at the subsolar point the average positions of the magnetopause and the shock wave are at  $\sim 10 R_E$  and  $13.5 R_E$  and flare out to  $\sim 13 R_E$  and  $22 R_E$ , respectively, at  $\theta_{SE} = 270^\circ$ . The solid lines indicate the result of a theoretical model calculation of the shock position with the magnetospheric boundary used as the blunt body [Spreiter and Jones, 1963]. Calculations of the shape of the sunward magnetopause [cf. Beard, 1960; Midgely and Davis, 1962; Spreiter and Briggs, 1962a, b; Mead and Beard, 1964] are in substantial agreement with these observations.

Explorer 10 observations of the magnetopause toward the tail of the magnetosphere as determined from simultaneous plasma [Bonetti et al., 1963] and magnetic field [Heppner et al., 1963] measurements have been rotated into the ecliptic plane for comparison and strengthen the experimental evidence that no closure of the geomagnetic tail is discernible contrary to closures predicted due to transverse plasma pressures [cf. Johnson, 1960; Beard, 1960].

Motions of the magnetospheric boundary and of the shock wave are implicit in the variations of the measurements summarized in Figure 4 [Ness et al., 1964] and are manifestations of temporal variations in the solar wind densities, velocities, and associated interplanetary magnetic fields. Explorer 10 observations in the evening flank of the magnetosphere over the radial distance range of  $\sim 22 R_E$  to  $42 R_E$  of alternate periods of no plasma and magnetic field directed radially from the sun and of plasma moving generally from the direction of the sun and disordered magnetic fields were interpreted as the motion of the magnetopause across the satellite's position. Freeman [1964] has shown that the correlation of  $D_{ST}$  values with the position of the sunward magnetopause as determined by Explorer 12 observations indicates that the magnetosphere is compressed during the initial phase of a magnetic storm presumably by enhanced solar plasma flux and is characterized by a general expansion during the recovery

phase. A single Explorer 12 observation of a transient presence of trapped protons near the magnetopause has given an impulsive, outward radial velocity of the boundary position exceeding  $150 \text{ km (sec)}^{-1}$  [Konradi and Kaufmann, 1965].

Geomagnetic tail region. The form of the distant geomagnetic field in the anti-solar direction has been of theoretical interest before definitive measurements were available: [cf. Piddington, 1959; Beard, 1960; Axford and Hines, 1961]; active investigations concerning the linear dimensions of the geomagnetic tail in the anti-solar direction are presently continuing [Dessler, 1964; Dungey, 1965; Michel, 1965; Piddington, 1965]. Explorer 10 [Heppner et al., 1963] provided the first direct experimental evidence indicating the nature of the magnetic field topology in the tail of the magnetosphere. Beyond approximately  $8 R_E$  to  $22 R_E$  a magnetic field directed radially outward (the probe was at southerly latitudes during the measurements) from the sun and considerably larger in magnitude over the theoretical values was detected. Beyond  $22 R_E$  the repeated transversal of the magnetopause across the satellite's position occurred. Measurements again at southerly ecliptic latitudes and to a radial distance of  $\sim 16 R_E$  with Explorer 14 [Cahill, 1964] provided confirmation of the above Explorer 10 results and indicated a depression of the magnetic field below theoretical values at 7 to  $8 R_E$ . Similar depressions of



the geomagnetic field were observed during the period near solar maximum with Explorer 6 (August 1959) at altitudes as low as  $5 R_E$  [Sonett et al., 1960 a, b; see also Smith, 1963, for a review of Pioneers 1 and 5, Luniks 1 and 2, and Explorers 6 and 10 measurements of the distant magnetic field and the possibility of an associated terrestrial ring current]. The most comprehensive measurements of the magnetic fields in the magnetospheric tail region reported thus far have been obtained by Ness [1965] with IMP 1 (apogee  $31.4 R_E$ ). A representative inbound pass [Ness, 1965] through the tail region is reproduced in Figure 5; the magnitude  $\overline{F}$  and the direction in solar ecliptic coordinates of the measured field are compared with the usual theoretical field (dashed lines). Ness concludes that (1) the large magnetic field of 10 to 30 gammas and characterized by regularity in direction observed to apogee distances ( $\sim$  half the distance to the lunar orbit) indicate that the earth's magnetic tail extends to at least the lunar orbit and (2) the simultaneous observations of a practically null field and a reversal from anti-solar direction to solar direction of the field at  $\sim 16 R_E$  is indicative (supported by similar observations over 18 consecutive orbits) of a neutral sheet in the tail region with a thickness of a fraction of an earth radius. Hence the earth's polar cap lines of force are pulled to the night side of the earth in the

magnetic tail with lines of force from the Northern and Southern Hemispheres directed along the solar and anti-solar directions, respectively, with these two flux tubes separated by a neutral sheet ( $B \simeq 0$ ) containing sufficient plasma to provide a pressure to balance the magnetic pressure of the adjacent tubes of force [see Axford et al., 1964, for a theoretical discussion]. A summary due to Ness [1965] of the magnetic field topology in the noon-midnight meridional plane is shown in Figure 6. Cahill [1965] has recently reported similar results near the midnight meridian which are summarized in Figure 7. The shape of the magnetic field has been projected into a magnetic meridian plane and significant stretching of the field lines in the anti-solar direction appears beyond  $\sim 10 R_E$ ; observation of a sharp field reversal indicative of the neutral sheet found by Ness was also reported. A coarse comparison [Obayashi, 1965] of the measured magnetic field magnitudes and the extrapolated field along the earth-to-sun line ( $X_{SE}$  axis) is given in Figure 8 and illustrates several of the experimental measurements discussed above. The magnetic field topology at high latitudes and large radial distance yet awaits experimental survey.

#### IV. OBSERVATIONS OF CHARGED PARTICLES IN THE DISTANT GEOMAGNETIC FIELD

The spatial distributions of charged particles surrounding and within the distant geomagnetic field are highly dependent upon the energy and species of charged particles under investigation. An attempt to summarize the major reported evidences provides insight into the complexity of the magnetospheric structure. Other summaries pertaining the the distributions of charged particles within the earth's magnetosphere have been given, for example, by Van Allen [1961], Farley [1963], O'Brien [1963], Frank and Van Allen [1964a], and Obayashi [1965].

Low-energy charged particles. An example of observations of a quasi-thermalized plasma just beyond the magnetopause and compressing the sunward geomagnetic field has been given previously in Figure 2 [Freeman et al., 1963]. The plasma energy flux was of the order of tens of ergs  $(\text{cm}^2\text{-sec-sr})^{-1}$  and was inferred from rudimentary dynamical considerations to be electrons of energy  $\sim$  few kilovolts and intensities  $J_{\ominus} \sim 10^9 - 10^{10} (\text{cm}^2\text{-sec})^{-1}$ . Measurements of the proton and electron intensities in this portion of the magnetosphere were continued with a plasma cup on IMP 1 by Bridge and his associates [1964, 1965] and a typical pass is reproduced here in Figure 9. Shown are the average proton energy  $V$ ,  $\log_{10}$  of the total positive ion flux and of the electron flux 65-210 eV looking toward and away from the sun and a measure,  $T$ , of the width of the proton distribution. Beyond  $15 R_E$

the satellite is in the interplanetary medium and a anisotropic flux of protons  $J_0 \simeq 10^8 \text{ (cm}^2\text{-sec)}^{-1}$  is directed away from the sun; the electron intensity 65 eV to 210 eV was reported to be at approximately the noise level of the instrument,  $\sim 6 \times 10^6 \text{ (cm}^2\text{-sec)}^{-1}$ . In the transition region between  $\sim 10 R_E$  and  $15 R_E$  approximately isotropic but variable intensities of protons and electrons  $J_0 \simeq 10^9 \text{ (cm}^2\text{-sec)}^{-1}$ , a marked increase in ion temperature (measurable intensities in all channels 45 eV through 5400 eV), and a sharp decrease in the intensities of protons and electrons at the magnetopause near  $10.5 R_E$  were observed. These measurements indicate that the proton density in the transition region exceeds the solar wind proton density by a factor of 8 to 10. A summary of the portions of the IMP 1 trajectory along which a hot, approximately isotropic proton flux was observed is shown in Figure 10 [Bridge et al, 1964, 1965]; the transition regions as determined by these plasma observations and by simultaneous magnetic field measurements are coincident (cf. Figure 4 [Ness et al., 1964] of this paper). Wolfe and Silva [1964] generally confirm the above results using a cylindrical analyzer also on IMP 1 and find that the proton angular distributions in the transition region indicate a highly turbulent flow statistically directed away from the sun. Measurements of low-energy electron intensities utilizing a retarding potential analyzer (also IMP 1, one outbound pass) have been obtained by Serbu [1965] and are shown here in Figure 11. The

omnidirectional intensity of electrons  $E > 100$  eV is  $\sim 2 \times 10^8$   $(\text{cm}^2\text{-sec})^{-1}$  from  $8 R_E$  to the position of the shock front at  $16 R_E$  as determined by the magnetometer measurements. Intensities of these electrons decrease sharply at the shock boundary but no change in the observed intensities, increase or decrease, is seen at the magnetopause at  $\sim 11 R_E$ . The radial profiles of the omnidirectional intensities of electrons  $0 < E < 5$  eV and  $5 < E < 10$  eV appear to enjoy a similar independence of the magnetospheric boundary.

A summary of the observations of the low-energy charged particle distributions near the sunward magnetopause as discussed above and in the tail of the magnetosphere is given in Figure 12. The ecliptic plane has been selected here for mapping and comparing the results. Explorer 10 plasma observations [Bonetti et al., 1964] indicate that the plasma resumes supersonic flow along the flanks of the magnetosphere. Observations of low-energy electron  $E \gtrsim 100$  eV intensities by Lunik 2 [Gringauz et al., 1961], Explorer 12 [Freeman, 1963], and Mars 1 [Gringauz et al., 1964] beyond  $\sim 8 R_E$  are in coarse agreement although there is an obvious need for more definitive measurements. The results of Gringauz and his collaborators [1964] at high latitudes ( $\lambda_m \sim 50^\circ$ ) (for coordinates, see Vakulov et al., [1964]) have been coarsely translated into the ecliptic plane by direct L-shell mapping. The position of these large electron intensities agrees with the onset of change in character of the geomagnetic field

in the anti-solar direction measured with IMP 1 and Explorer 14 as reviewed above. The observed electron energy flux of several tens of ergs  $(\text{cm}^2\text{-sec})^{-1}$  (energy density  $\sim 10^{-8}$  ergs  $(\text{cm})^{-3}$  if  $E \sim$  a few kilovolts) appears sufficient to support the neutral sheet discovered by Ness [1965] if  $B \sim 30$  gammas for the magnitude of the adjacent fields. Large intensities of electrons of similar energies at low altitudes and high latitudes  $L \gtrsim 8$  have been observed during local night [Sharp et al., 1965; Fritz and Garnett, 1965] and further suggest the tail region as at least a partial reservoir and/or accelerating region for auroral particles. The spatial connection of the two regions characterized by large intensities of electrons  $E > 100$  eV near the sunward magnetopause and in the geomagnetic tail has not been determined [cf. Gringauz, 1964; Van Allen, 1964] and further surveys in the dawn, evening, and tail portions of the magnetosphere are necessary to establish the relationship.

Medium-energy charged particles. Simultaneous measurements of magnetic fields and electron  $E \sim 50$  keV intensities with Explorer 12 revealed that a large decrease of electron intensities occurred concurrently with the disappearance at the sunward magnetopause of the uniform, although perturbed, geomagnetic field necessary for the durable trapping of charged particles [Freeman et al., 1963]

(refer to present Figure 2). Another example of the typical behavior of the radial profiles of electron  $E > 40$  keV,  $E > 230$  keV, and  $E > 1.6$  MeV intensities near the sunward magnetospheric boundary with Explorer 14 [Frank and Van Allen, 1964b] is given in the left-hand inset of Figure 13. The "spikes" of electron  $E > 40$  keV intensities beyond the position of the trapping boundary at  $\sim 70,000$  km and extending in this example to  $\sim 93,000$  km in the general vicinity of the transition region have been observed by Anderson et al. [1965], Frank and Van Allen [1964b], and Fan et al. [1964] and are suggestive of a high-energy tail of the electron spectrums in the transition region ( $J_o(E > 40 \text{ keV})/J_o(E \gtrsim 1 \text{ keV}) \sim 10^{-5}$ ) [Frank and Van Allen, 1964b]. Theoretical hypotheses for the acceleration of these energetic electrons have been given by Scarf et al. [1965] and Jokipii and Davis [1964]. A coarse positive correlation of the occurrence of these intensity 'spikes' with the planetary magnetic index  $K_p$  (right-hand inset Figure 13) was found [Frank and Van Allen, 1964b] with average  $K_p$  values  $\simeq 2.5$  and  $\simeq 1.5$  for passes with and without intensity spikes beyond the magnetopause, respectively. The percentage of passes with observations of spikes while the satellite was in the transition region and  $K_p$  was 3 or 4 ( $\sim 20$  passes) was  $\sim 75\%$ . Figure 13 also displays the 'flaring' of the magnetospheric boundary as deduced

from the characteristic sharp termination of intensities in general agreement with IMP 1 magnetic field measurements [Ness et al., 1964] and the increasing thickness of the regions with extra-magnetospheric spikes with angular displacement from the subsolar point [Frank et al., 1963; Fan et al., 1964; Frank and Van Allen, 1964b; Anderson et al., 1965]. Direct correlation of the position of the shock as determined by simultaneous magnetic field measurements with the positions of these intensity spikes shows that the intensity spikes are not coincident with, but are also sometimes positioned upstream from, the position of the shock wave [Anderson et al., 1965; Jokipii and Davis, 1964]. The interpretation due to Anderson et al. [1965] of small intensity spikes ( $\sim 10^3$  to  $10^4$  (cm<sup>2</sup>-sec)<sup>-1</sup>) observed at large distances in front of the shock as being of terrestrial origin may be complicated by the recent discovery [Van Allen and Krimigis, 1965] of solar electron  $E \sim 40$  keV events with Mariner 4 data.

A summary of several gross features of the spatial distributions of electrons  $E \sim 40$  keV is given in Figure 14. A rough upper limit of the average electron  $E > 40$  keV intensity in the interplanetary medium is  $\sim 10$  (cm<sup>2</sup>-sec)<sup>-1</sup> as determined with Mariner 4 [cf. Van Allen, 1965]. Surrounding the earth and extending to the magnetopause in the solar direction ( $\sim 10 R_E$ ) and to  $\sim 8 R_E$  in the anti-solar direction is a 'hard core' of electron  $E > 40$  keV intensities



$\sim 10^7\text{-}10^8 \text{ (cm}^2\text{-sec)}^{-1}$  characterized by relatively weak temporal and spatial variations when compared to outlying regions of the magnetosphere [Frank, 1965a]. Within  $45^\circ$  of longitude of the subsolar point at low latitudes these intensities are terminated coincidentally with the magnetopause. In the dawn and evening 'skirts' of the spatial (i.e., at larger angular displacements from the subsolar point) distribution tortuous temporal and spatial variations are observed [Frank et al., 1963; Anderson et al., 1965; Frank, 1965a] and their association with the magnetopause in terms of spatially correlated decreases or increases in intensity weakens and disappears toward the flanks of the magnetosphere [Anderson et al., 1965]. Montgomery et al. [1965] have noted a large asymmetry in the dawn and evening skirts with the dawn skirt having generally more intense and frequent electron  $E > 50$  keV intensity spikes when compared with observations in the local evening skirt. Observations with Explorer 14 displayed no significant dawn-evening asymmetry during relatively quiescent magnetic conditions ( $\sum K_p \leq 25$ ) [Frank, 1965a] but the nearly circular Vela orbits [Montgomery et al., 1965] are more favorable for observing such an asymmetry than the highly eccentric Explorer 14 orbit. In the anti-solar direction a 'tail' of electron  $E > 40$  keV intensities has been found [Frank, 1965a; Anderson, 1965] characterized by much weaker temporal and spatial

dependences when compared to those of the skirt regions and a shallow latitude dependence compared to that of the 'core' of intensities. The position of this 'tail' of electron intensities presumably is coincident with the neutral sheet. Anderson [1965] has found 'spikes' of intensities with fast risetimes ( $\sim$  a few minutes) with approximately exponential decays ( $\sim$  a few minutes to hours) in the anti-solar direction to at least IMP 1 apogee position of  $\sim 32 R_E$  with a strong decrease in frequency of occurrence with increasing radial distance and infers that these 'islands' (typically several  $R_E$  in linear dimensions) are impulsively injected into the tail region. Asbridge et al. [1965] have observed electrons of similar energies streaming in the anti-solar direction in the geomagnetic tail at  $\sim 17 R_E$ . Recently intensity 'spikes' of low-energy protons  $E > 125$  keV,  $j \simeq 10^3$  to  $10^4$  ( $\text{cm}^2\text{-sec-sr}$ ) $^{-1}$  in these regions have been reported [Konradi, 1965]. Observations of rapid increases of intensities during the onsets of several intensity 'spikes' in the evening skirt of intensities with OGO 1 allowed Frank et al. [1965] to set a lower limit on the bulk velocity of the irregularities of  $50 \text{ km (sec)}^{-1}$ . A coarse summary of the average omnidirectional intensities of electrons  $E > 40$  keV as a function of geocentric radial distance in the anti-solar direction is shown in Figure 15 which includes the Explorer 14 and IMP 1 data

discussed above and a single determination at  $3300 R_E$  with Mariner 4 [Van Allen, 1965]; average intensities at the lunar orbit in the anti-solar direction might reasonably be expected to be  $\sim 10^3 \text{ (cm}^2\text{-sec)}^{-1}$ .

Periodic variations of the intensities of electrons  $E > 40 \text{ keV}$  with periods  $\sim 6$  minutes and persisting for several hours in the skirt and tail regions of the magnetosphere have been recently reported [Anderson, 1965]. These variations are presumably related to periodic variations in the magnetic field (cyclic acceleration and deceleration by hydromagnetic waves) such as the long-period hydromagnetic waves observed by Patel [1965] (for theoretical discussions of the generation and propagation of magnetospheric hydromagnetic waves, refer to Dessler [1958, 1961], Dessler and Parker [1959], Dessler and Walters [1964], and MacDonald [1961]).

Energetic electrons. The spatial distributions of electron  $E \gtrsim 1.5 \text{ MeV}$  intensities in the outer magnetosphere and beyond are summarized in Figure 16. The isopleths of constant intensities  $J_O = 10^2, 10^3, 10^4 \text{ (cm}^2\text{-sec)}^{-1}$  as measured with Explorer 14 near the magnetic equatorial plane exhibit a large local time dependence in radial position (for example, the isopleth  $J_O = 10^2 \text{ (cm}^2\text{-sec)}^{-1}$  is positioned at  $10 R_E$  in the solar direction and at  $8 R_E$  in the anti-solar direction) and the spatial distribution of these energetic electrons lies within the relatively

stable 'core' of 40 keV electron intensities (see Figure 14) [Frank, 1965b]; this region as defined by the outermost contour is probably indicative of the region in which complete longitudinal drift paths around the earth are possible. Several intensity spikes of electrons  $E > 1.5$  MeV ( $J_0 \sim 10 \text{ (cm}^2\text{-sec)}^{-1}$ ) were observed in the dawn skirt and transition region and none in the corresponding evening portions of the magnetosphere but the statistics are poor with only several events of observable intensity discovered in the data. Arrows on the isopleths indicate excursions (temporal variations) in the position of the contour  $J_0 = 10^2 \text{ (cm}^2\text{-sec)}^{-1}$  observed with Explorer 14. Upper limits on the average interplanetary omnidirectional intensities are  $\overline{J_0} (> 3 \text{ MeV}) \lesssim 0.1 \text{ (cm}^2\text{-sec)}^{-1}$  [Cline et al., 1964] and  $\overline{J_0} (> 1.5 \text{ MeV}) \lesssim 1 \text{ (cm}^2\text{-sec)}^{-1}$  [Frank, private communication] with IMP 1 and Explorer 14 data, respectively.

The diurnal variation of electron  $E > 280$  keV intensities observed at low altitudes  $L \gtrsim 6$  [Williams and Palmer, 1965] is in qualitative agreement with the observations of electrons  $E > 1.5$  MeV in the magnetic equatorial plane [Frank, 1965b] and is consistent with the motion of charged particles in models of the perturbed geomagnetic field [cf. Hones, 1963].

Determinations of the isopleths of electron  $E > 1.5$  MeV intensities [Frank, 1965b] in the outer magnetosphere indicate

that the axis of symmetry of the local time distribution of charged particles may not be the earth-to-sun direction but may lie closer to the 11:00-23:00 (local time) meridional plane. Although this single result is inconclusive, a similar effect has been observed in the diurnal variations of electron  $E > 40$  keV intensities at low altitudes with Injun 3 data [Frank et al., 1964].

## ACKNOWLEDGEMENTS

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## FIGURE CAPTIONS

Figure 1. Illustrations of the solar ecliptic coordinate system (upper diagram) and of the solar magnetospheric coordinate system (lower diagram).

Figure 2. Particle and magnetic field measurements with Explorer XII for the inbound pass on 13 September 1961. The CdSB (magnet in aperture) detector count-rate has been normalized to the energy scale of the CdSTE (no magnet in aperture) detector. The counting rates of both CdS detectors are nearly linear with energy flux. Both spectrometer channels (SpL and SpH) have been corrected for background counts by the subtraction of the counting rate of the background detector SpB. The CdS optical monitor (not shown) indicated that during this pass the CdS detectors did not have any bright objects within their field of view.  $F$  denotes the scalar magnetic field strength;  $\alpha$  the angle between the  $F$ -vector and the spin-axis of the satellite; and  $\psi$  the dihedral angle between the plane containing the  $F$ -vector and the spin axis and the plane containing the spin axis and the satellite-sun line (after Freeman et al. [1963]).

Figure 3. Magnetic field data (IMP 1) from orbit No. 11 January 5, 1964, illustrating outbound traversal of magnetopause boundary at  $13.6 R_E$  and shock wave boundary at  $19.7 R_E$  (after Ness et al. [1964]).

Figure 4. Summary of location of rectified shock wave and magnetopause boundary crossings and comparison with theoretical predictions using analogy with supersonic fluid flow (after Ness et al. [1964]).

Figure 5. Magnetic field data from inbound orbit No. 41, May 2 through May 4, 1964. Throughout most of this time interval the magnetic field is pointed anti-solar although at a geocentric distance of  $16 R_E$  the magnetic field abruptly reverses direction at the same time that it becomes very weak or zero. This spatially limited region is identified as a neutral sheet in the earth's magnetic tail (after Ness [1965]).

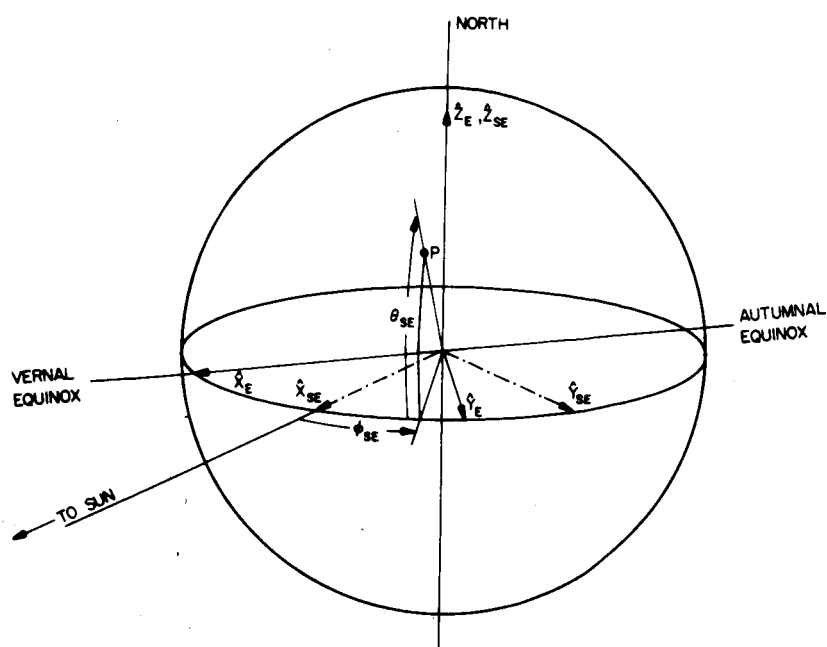
Figure 6. Illustration of the interpretation of the IMP 1 magnetic field data field topology within the noon-midnight meridian plane. The relative position of the neutral surface or sheet in the earth's magnetic tail and the corotating magnetic field lines supporting trapped particle motion are indicated. These include the classical Van Allen radiation belts. The collisionless shock and magnetosphere boundaries are extrapolated to the polar regions indicating a depression but the relative position of a polar neutral point and the size of the boundary are not experimentally determined. Cylindrical symmetry about the earth-sun line has been assumed for the boundary of the earth's magnetic tail in this presentation (after Ness [1965]).



- Figure 7. A schematic representation of the magnetic lines of force projected upon a magnetic meridional plane as observed by Explorer 14 near local midnight (after Cahill [1965]).
- Figure 8. Schematic representation of the outer geomagnetic field in the equatorial plane along the sun-earth line (after Obayashi [1965]).
- Figure 9. Measurements of low-energy protons and electrons in the transition region and in the interplanetary medium with a Faraday cup type instrument on IMP 1 (after Bridge et al. [1965]).
- Figure 10. Summary of portions of IMP 1 trajectory during which a hot, nearly isotropic plasma flux was observed (after Bridge et al. [1965]).
- Figure 11. Plot of low-energy electron fluxes as a function of geocentric radial distance (after Serbu [1965]).
- Figure 12. Coarse summary of low-energy charged particle observations. Average positions of the magnetopause and the shock as given by Ness [1965] have been included in the diagram.
- Figure 13. Summary of Explorer 14 observations of the sharp termination of electron  $E > 40$  keV intensities at the sunward magnetopause and of the intensity 'spikes' in and near the transition region (after Frank and Van Allen [1964b]).
- Figure 14. Coarse summary of the observations of electrons  $E > 40$  keV near the ecliptic plane.

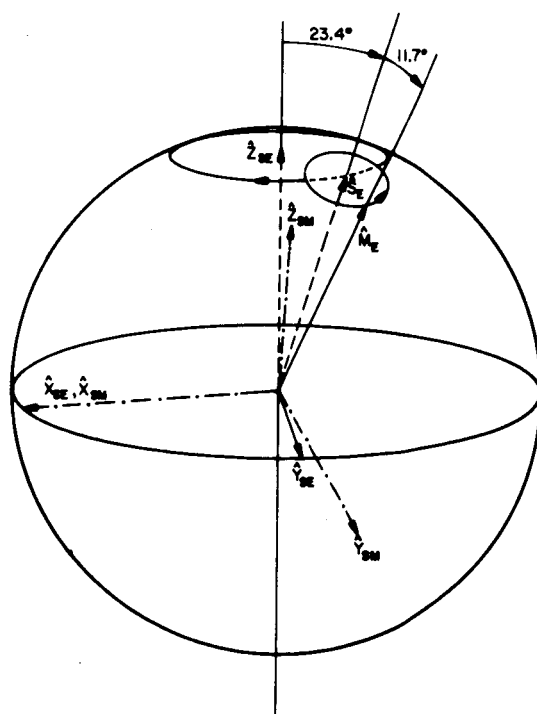
Figure 15. Summary of the observations of electron  $E > 40$  keV intensities as a function of radial distance in the anti-solar direction.

Figure 16. Coarse summary of the observations of electron  $E > 1.6$  MeV intensities in the magnetic equatorial plane.



SOLAR ECLIPTIC COORDINATE SYSTEM

$\hat{x}_E, \hat{y}_E, \hat{z}_E$  ECLIPTIC  
 $\hat{x}_{SE}, \hat{y}_{SE}, \hat{z}_{SE}$  SOLAR ECLIPTIC  
 $\theta_{SE}$  SOLAR ECLIPTIC LATITUDE  
 $\phi_{SE}$  SOLAR ECLIPTIC LONGITUDE



SOLAR MAGNETOSPHERIC COORDINATE SYSTEM

$\hat{M}_E$  || EARTH'S MAGNETIC  
 DIPOLE AXIS  
 $\hat{S}_E$  || EARTH'S SPIN AXIS  
 $\hat{x}_{SM}, \hat{y}_{SM}, \hat{z}_{SM}$  SOLAR MAGNETOSPHERIC  
 $\hat{x}_{SM} || \hat{x}_{SE}$   
 $\hat{y}_{SM} || \hat{M}_E \times \hat{x}_{SM}$

Figure 1

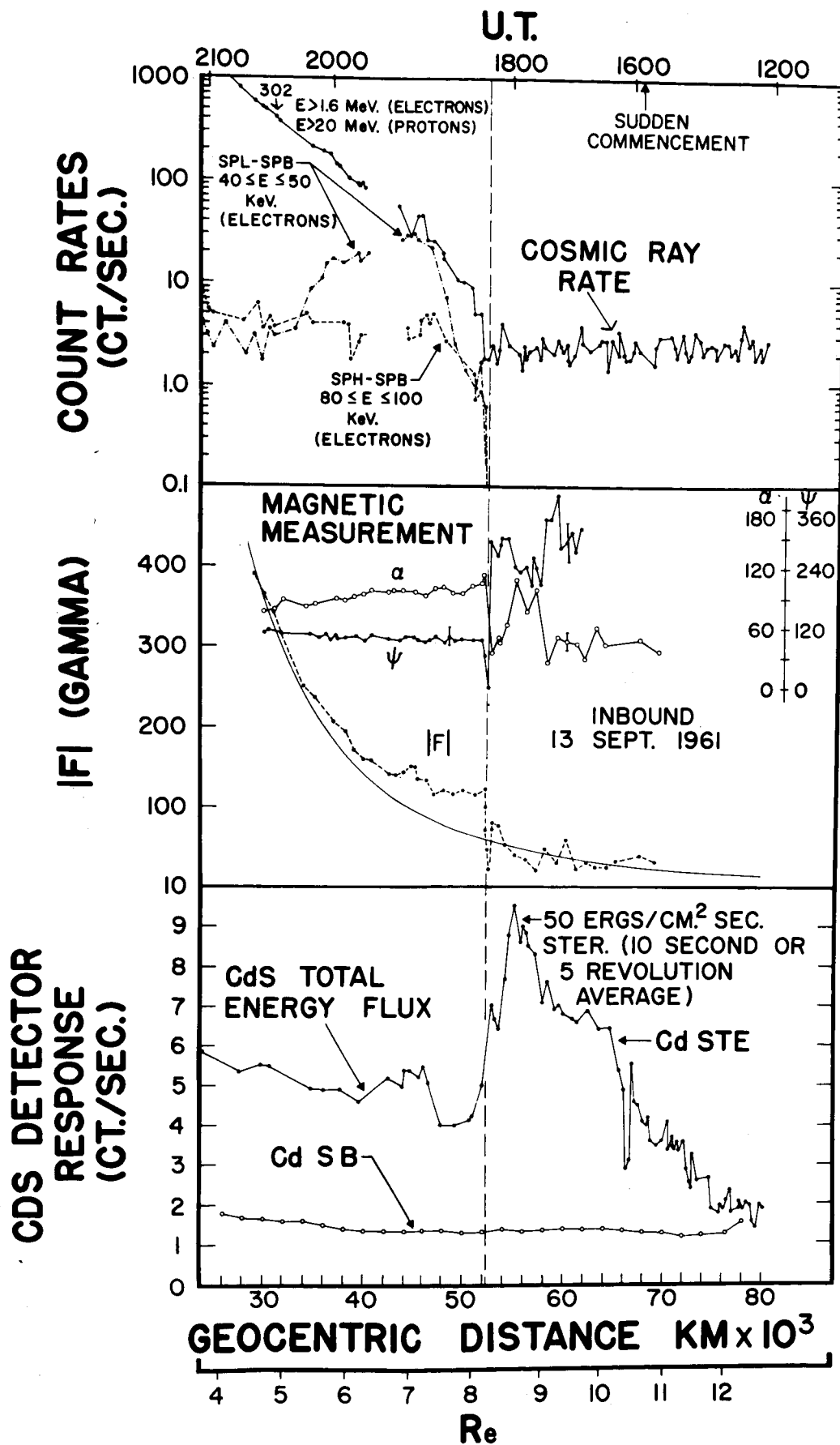


Figure 2

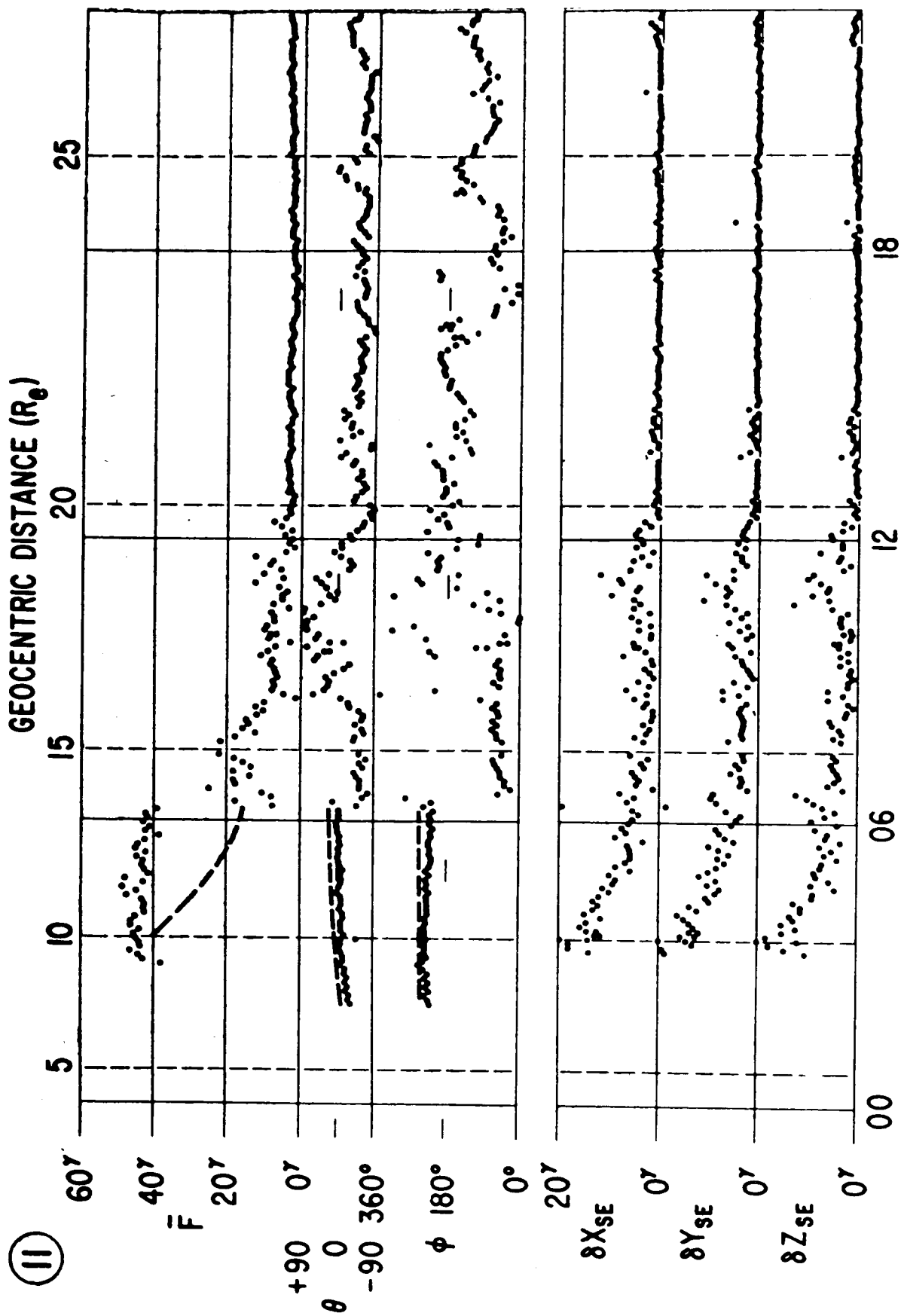


Figure 3

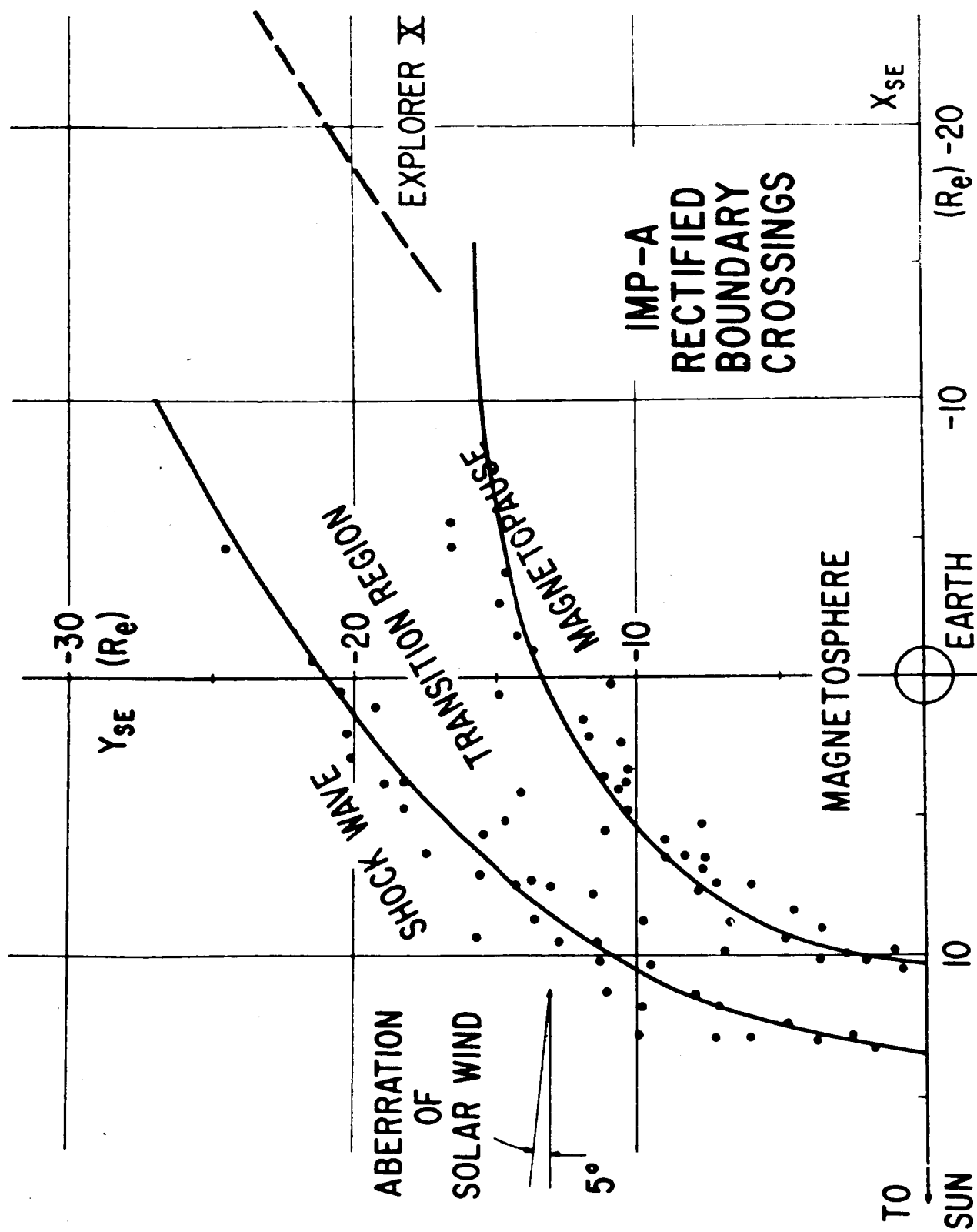
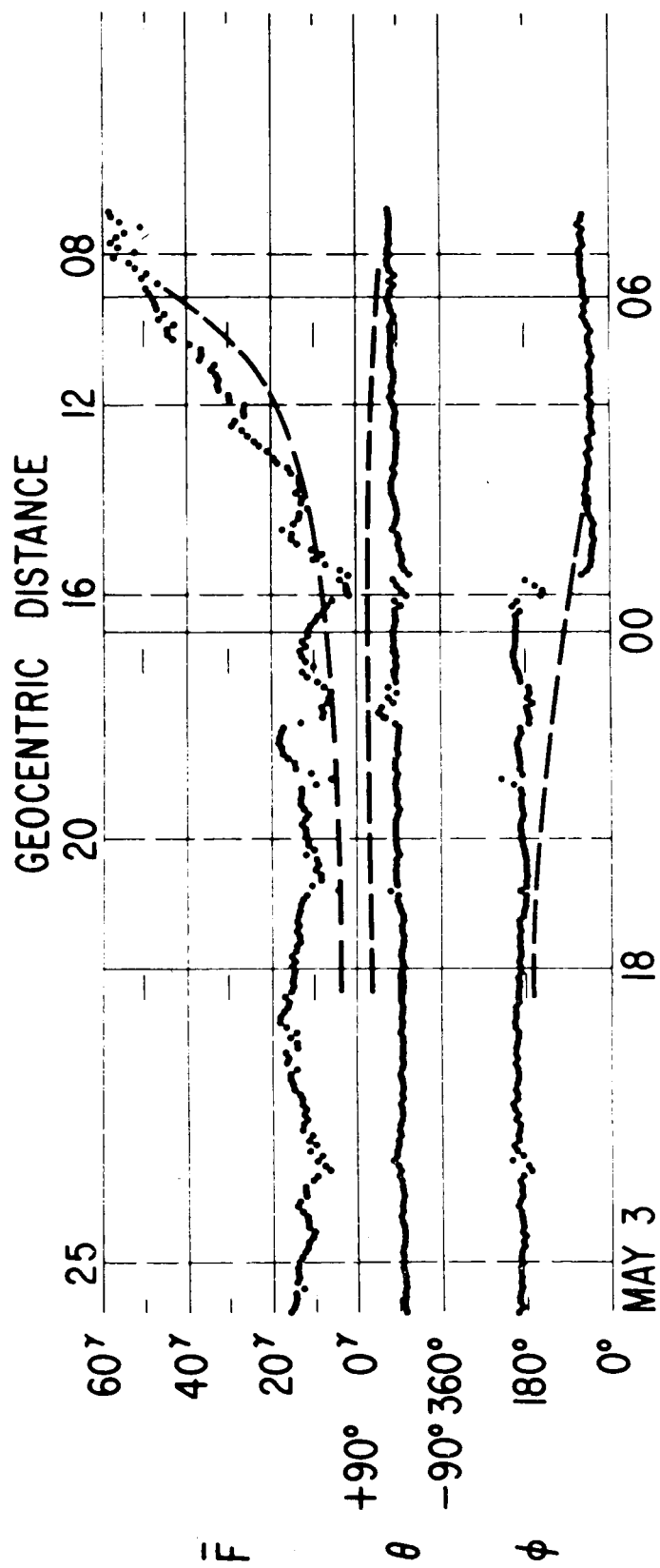
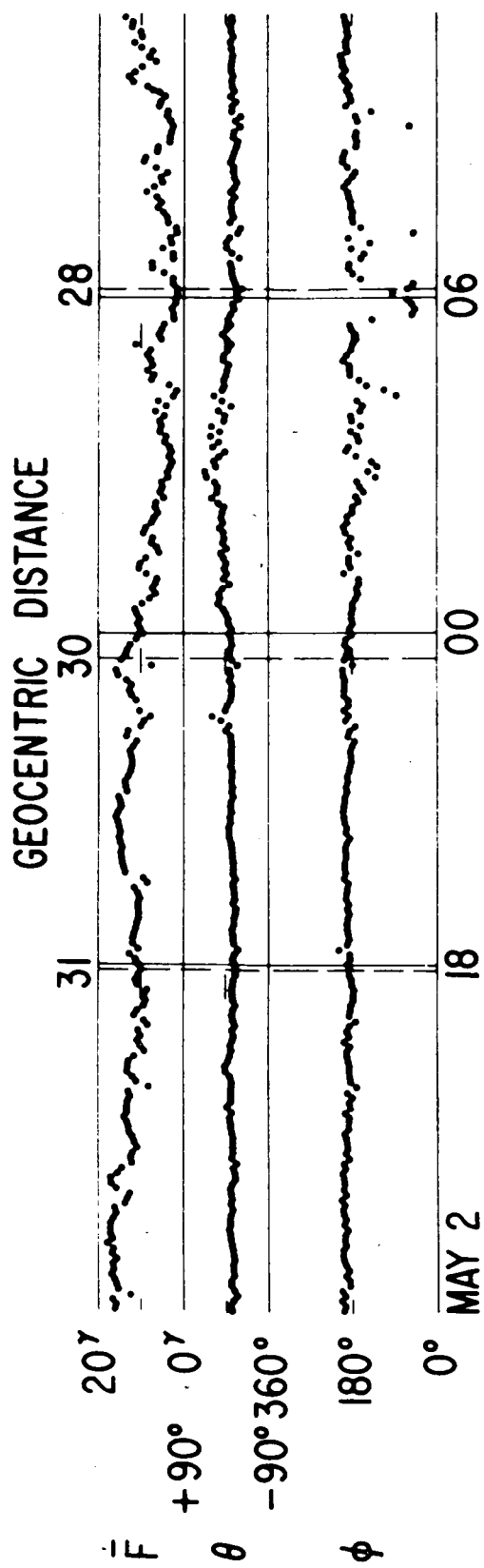


Figure 4



ORBIT NO. 41 IMP-1 1964

Figure 5

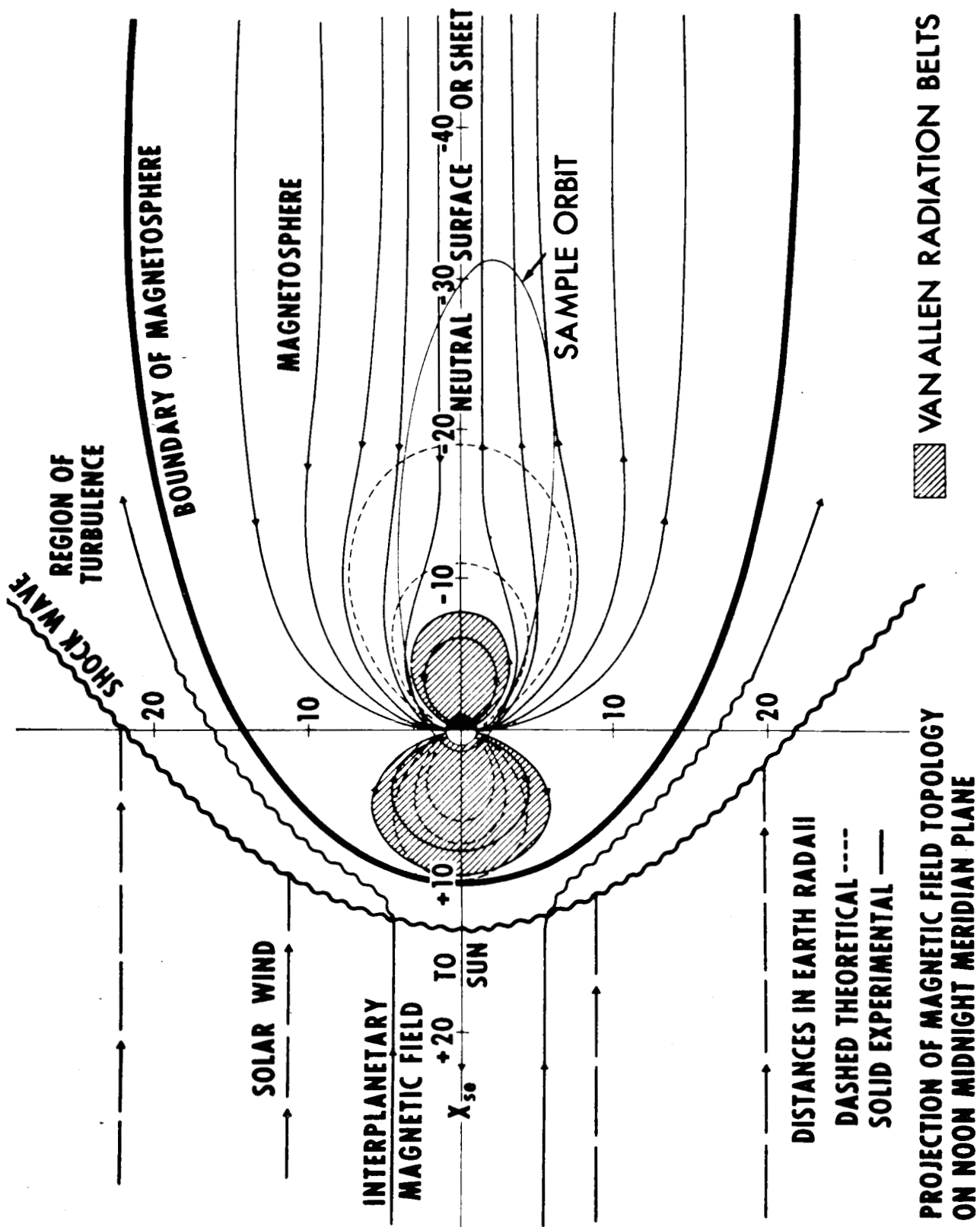


Figure 6



EXPLORER 14

JAN. 27 - FEB. 20, 1963

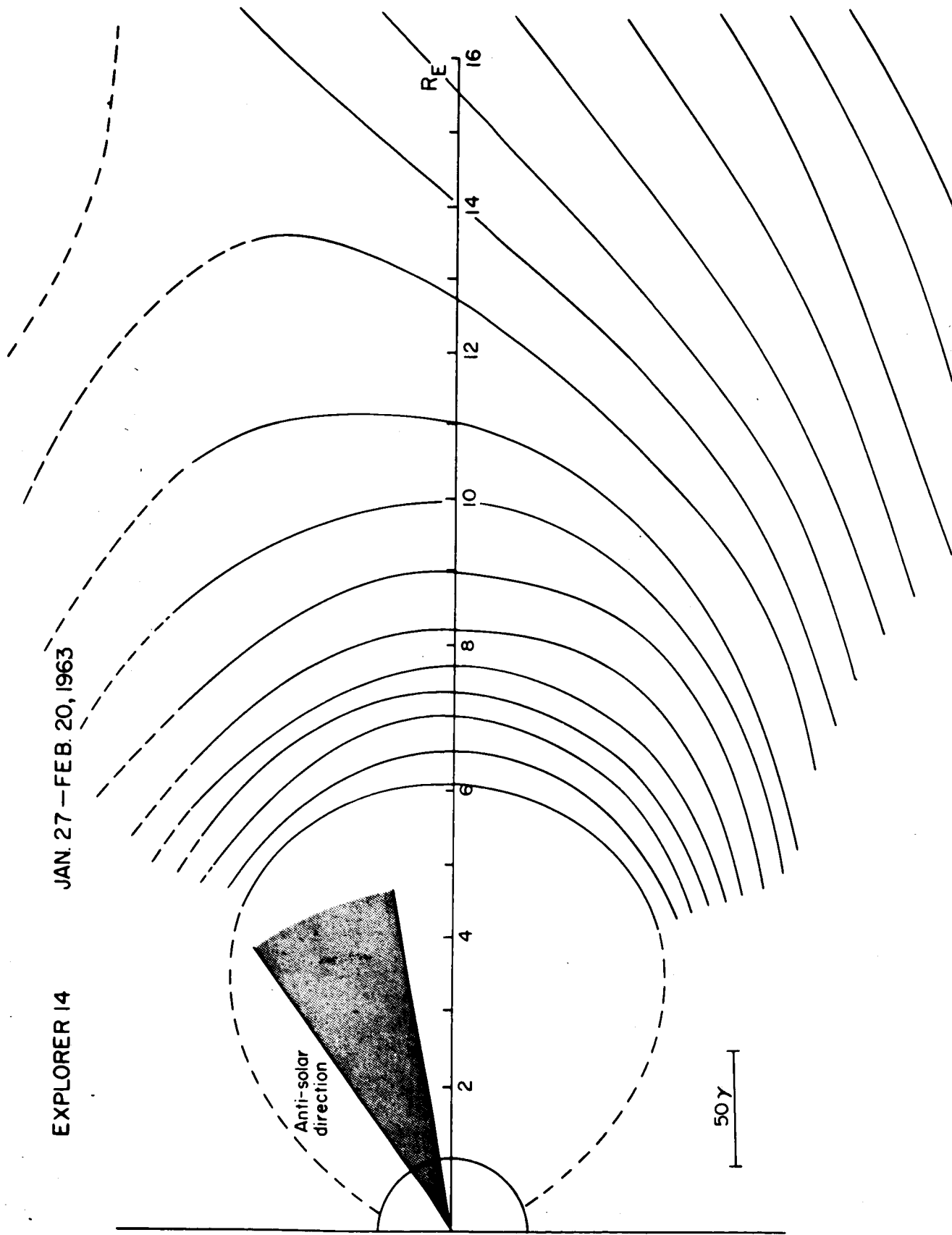
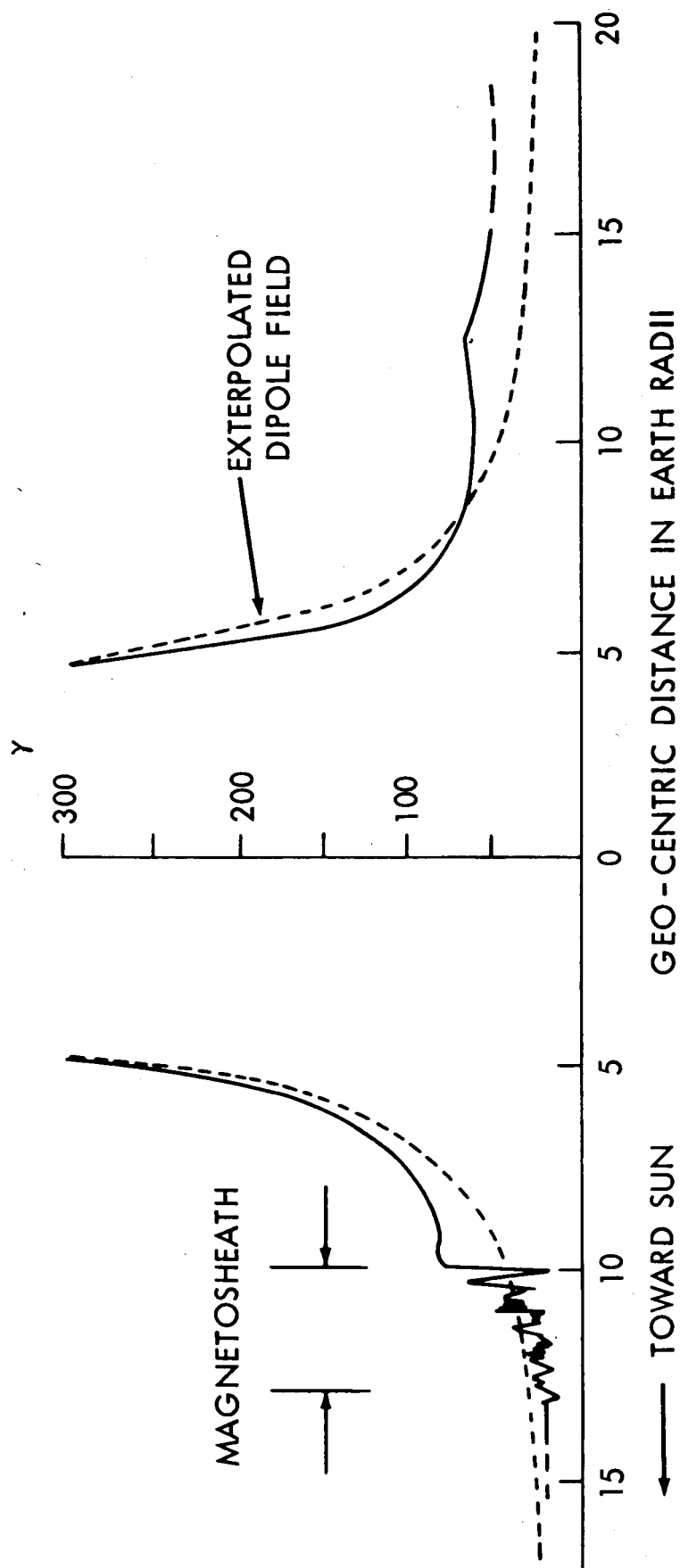


Figure 7



Schematic representation of the outer geomagnetic field in the equatorial plane along the sun-earth line.

Figure 8

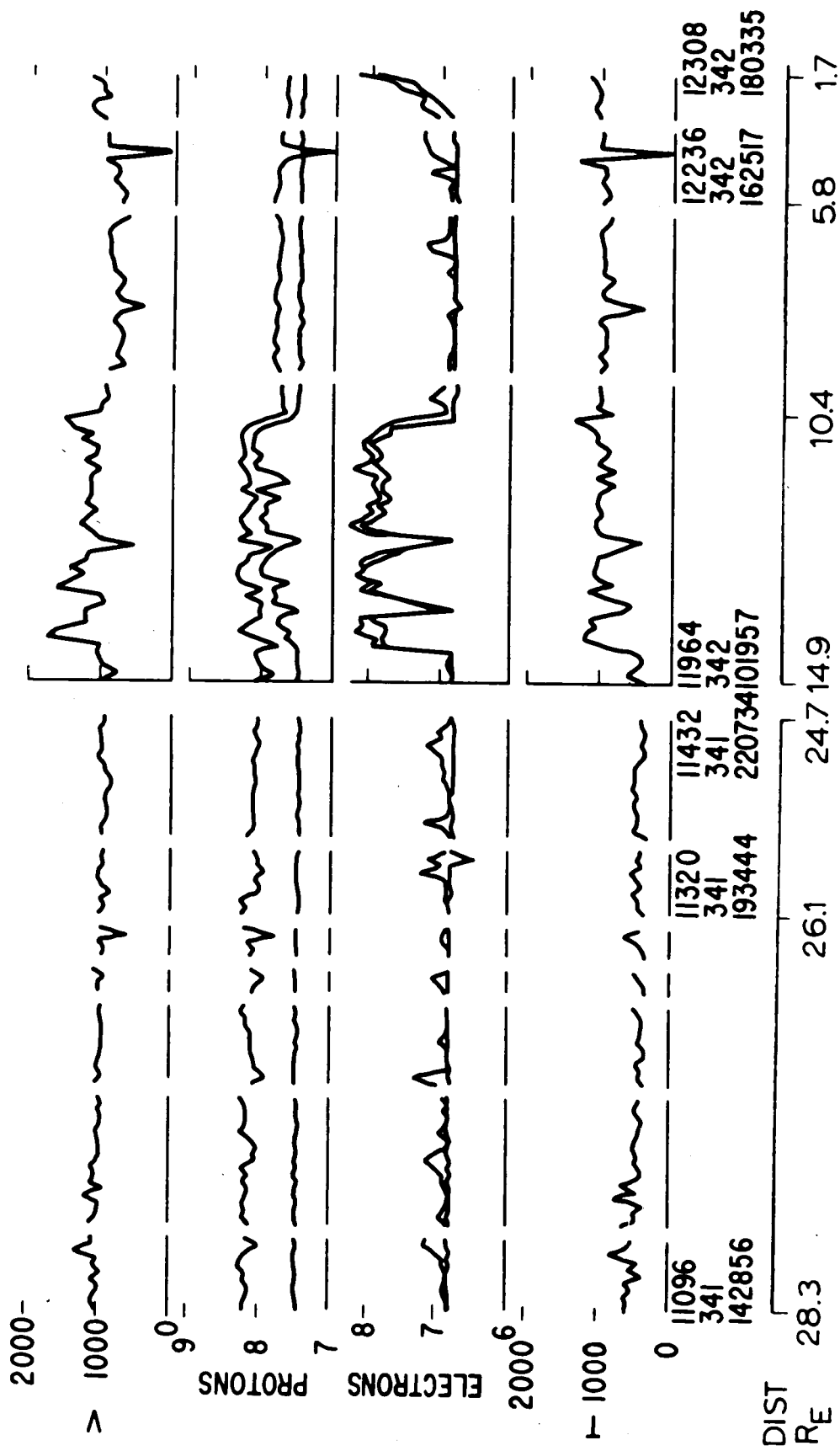


Figure 9

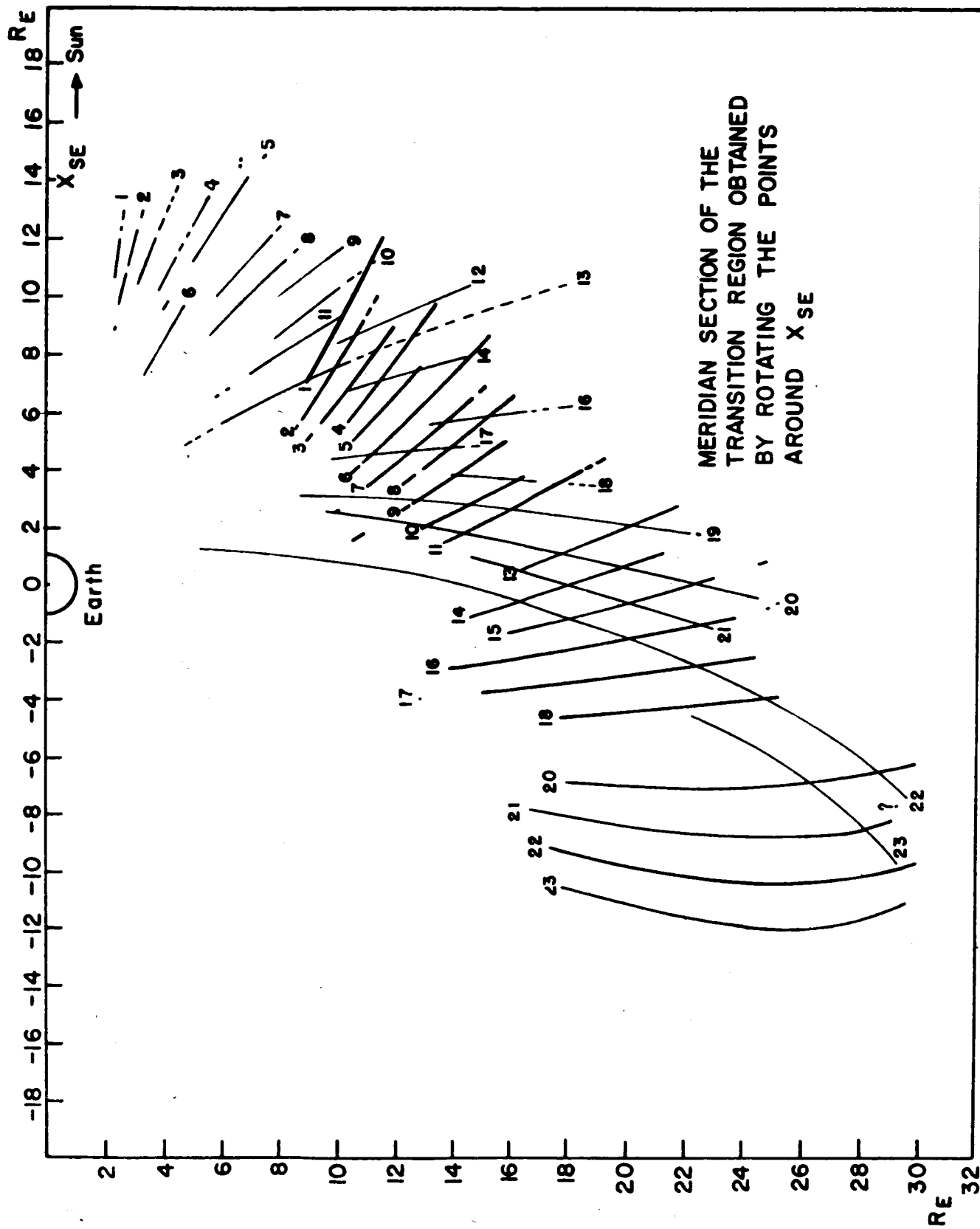


Figure 10

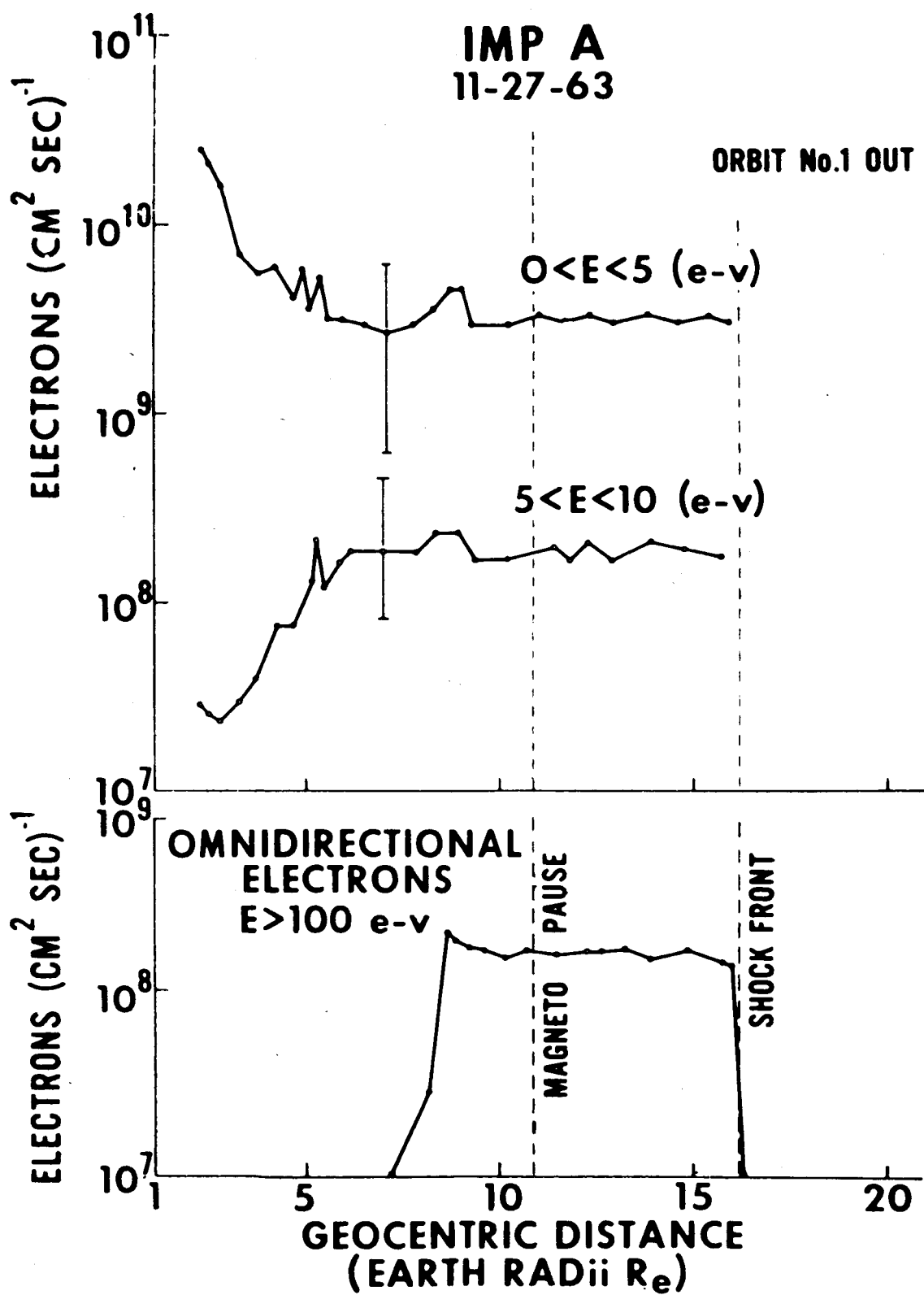


Figure 11

# PROTONS AND ELECTRONS $E \sim 1$ KeV

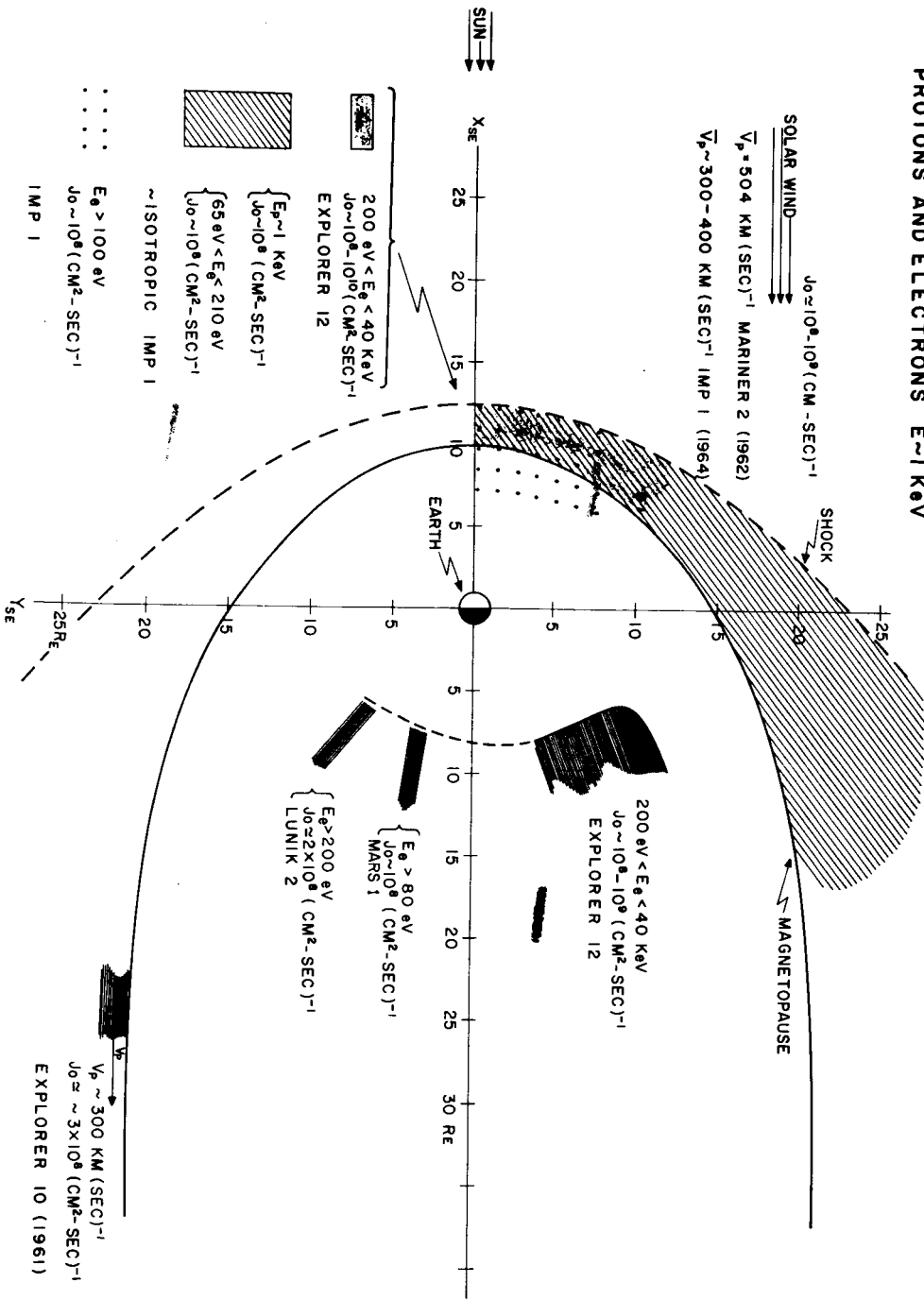


Figure 12

# PROTONS AND ELECTRONS $E \sim 1 \text{ KeV}$

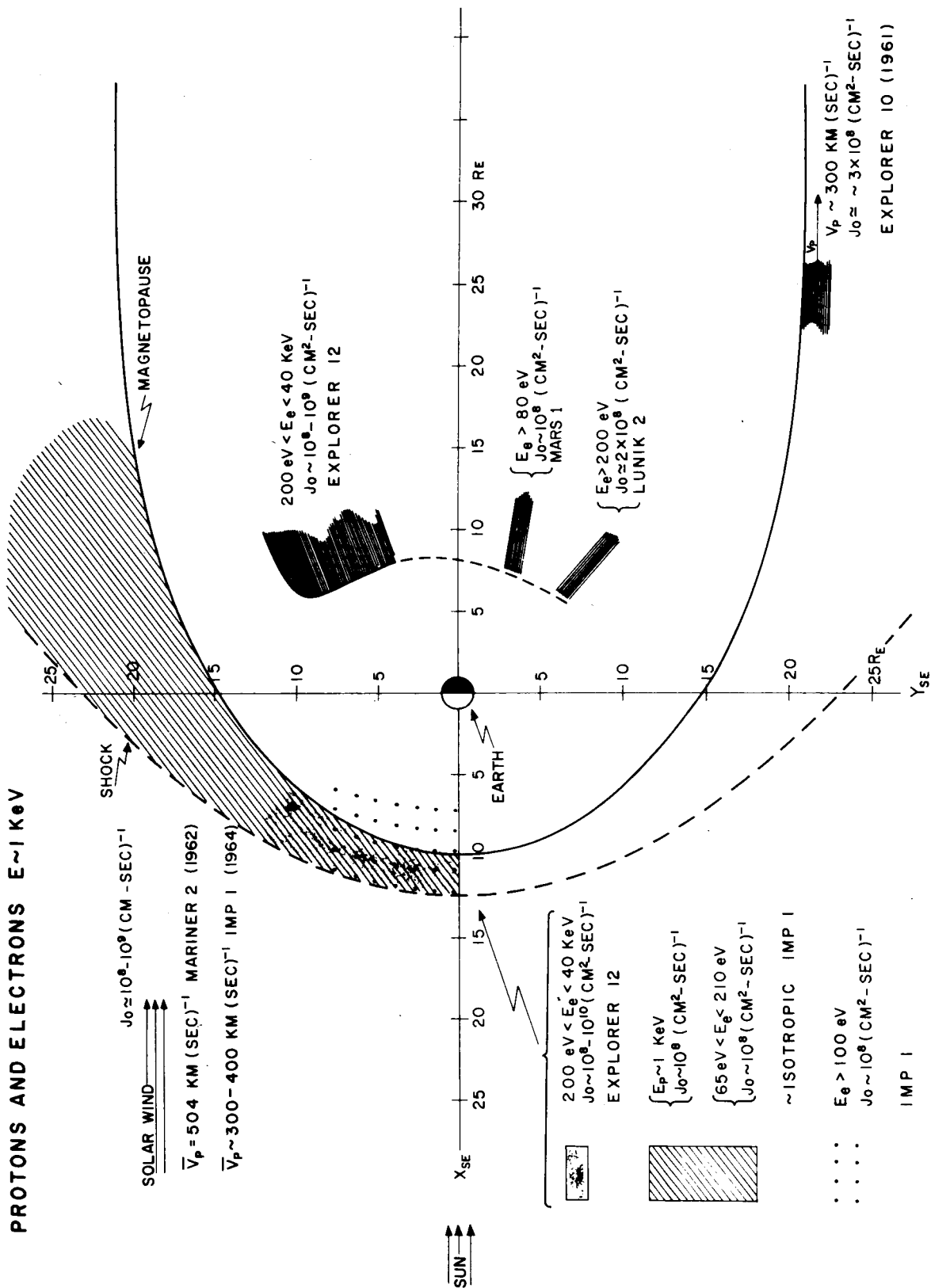
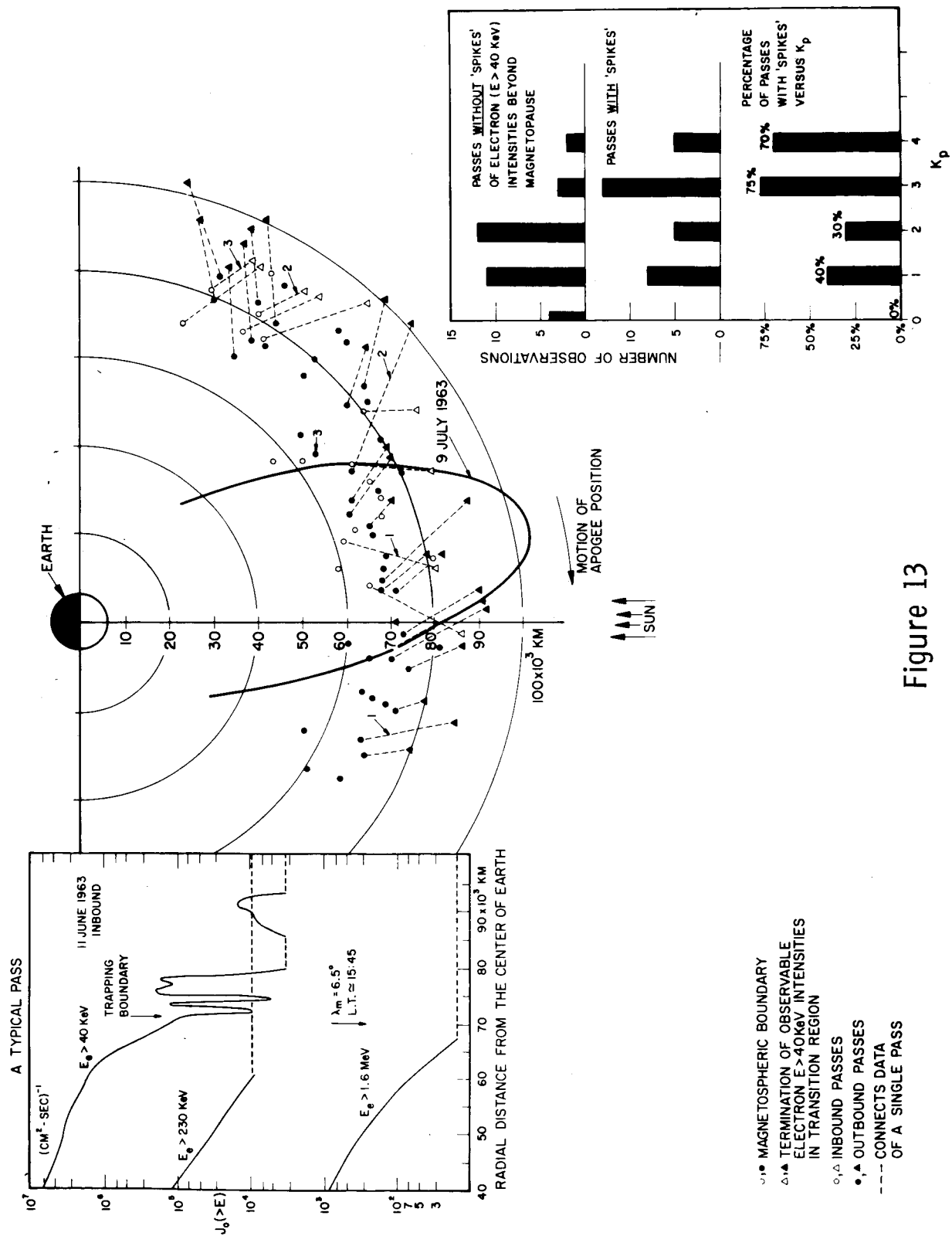


Figure 12

POSITIONS OF THE GEOMAGNETIC TRAPPING BOUNDARY AND OF THE  
TERMINATION OF ENERGETIC ELECTRONS IN THE TRANSITION REGION  
AS DETERMINED WITH EXPLORER XIV MEASUREMENTS OF ELECTRONS  $E > 40 \text{ KeV}$

**MAY - AUGUST, 1963**



## Figure 13



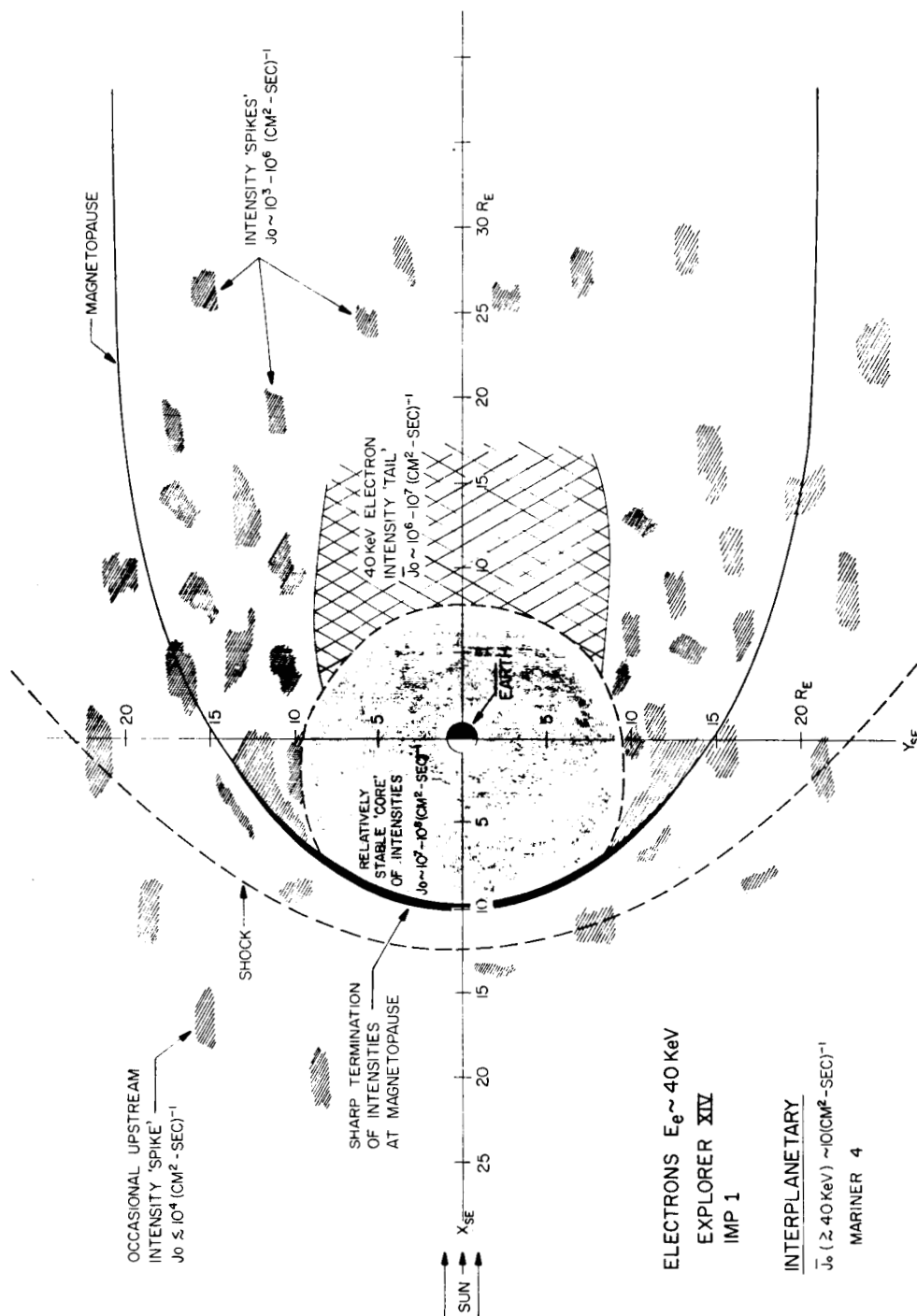
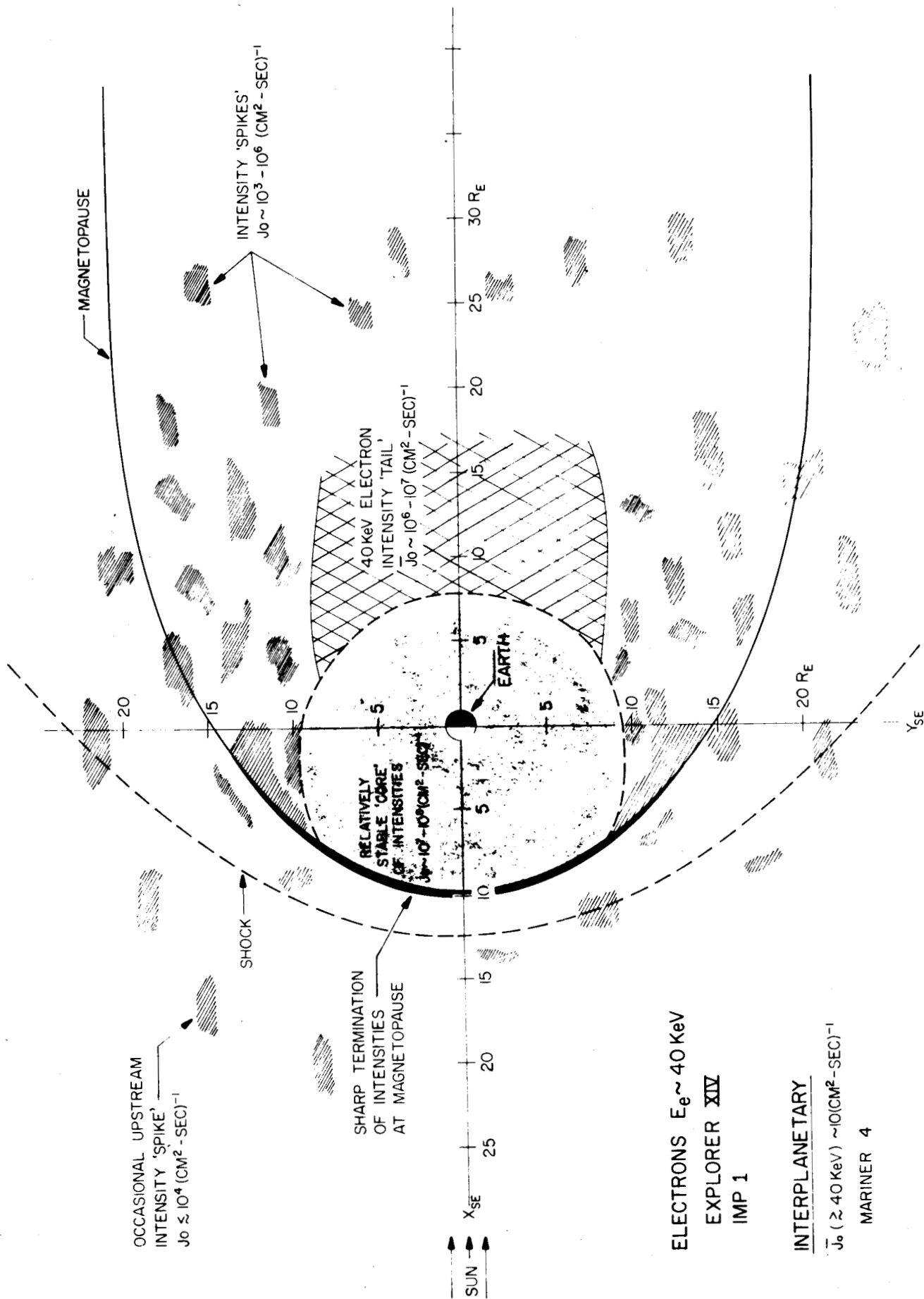


Figure 14



ELECTRONS  $E_e \sim 40 \text{ KeV}$   
EXPLORER XIV  
IMP 1

INTERPLANETARY

$J_0 (\geq 40 \text{ KeV}) \sim 10 \text{ (CM}^2\text{-SEC)}^{-1}$

MARINER 4

Figure 14

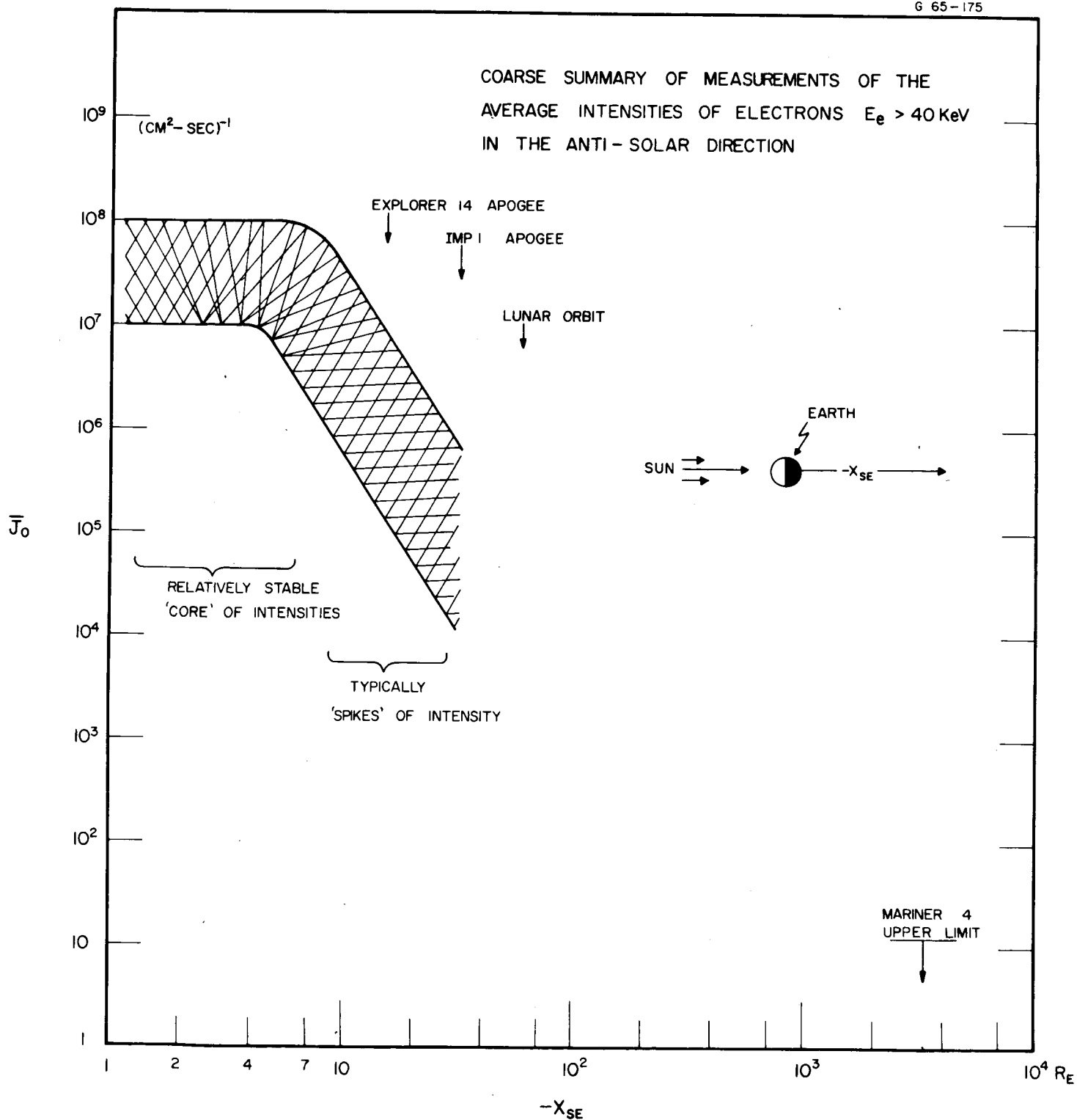


Figure 15

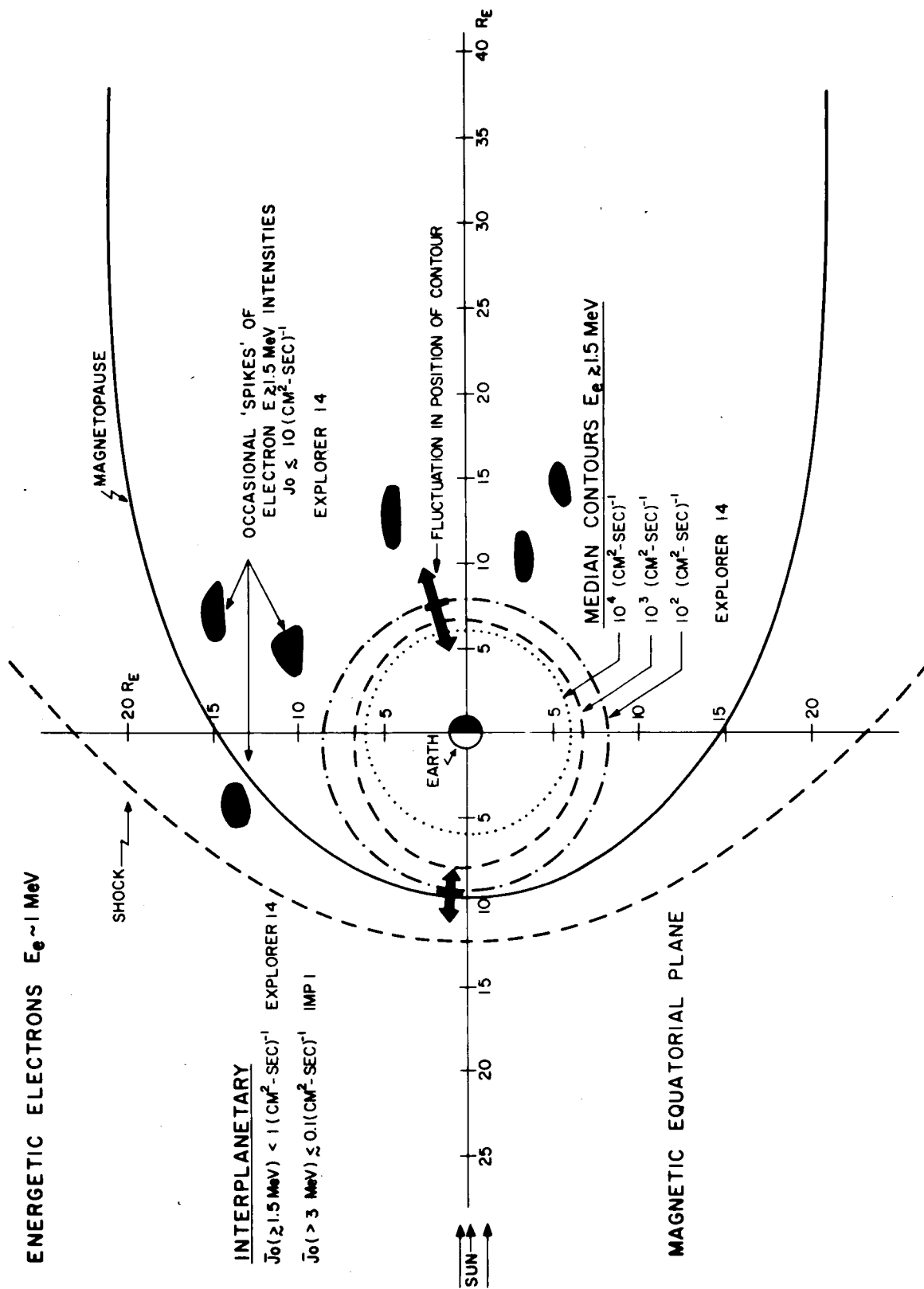


Figure 16

UNCLASSIFIED

Security Classification

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1. ORIGINATING ACTIVITY (Corporate author) University of Iowa, Department of Physics and Astronomy		2a. REPORT SECURITY CLASSIFICATION UNCLASSIFIED	
		2b. GROUP	
3. REPORT TITLE Observations of Magnetospheric Boundary Phenomena			
4. DESCRIPTIVE NOTES (Type of report and inclusive dates) Progress August 1965			
5. AUTHOR(S) (Last name, first name, initial) Frank, L. A.			
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8a. CONTRACT OR GRANT NO. Nonr-1509(06)		9a. ORIGINATOR'S REPORT NUMBER(S) U. of Iowa 65-32	
b. PROJECT NO.			
c.		9b. OTHER REPORT NO(S) (Any other numbers that may be assigned this report)	
d.			
10. AVAILABILITY/LIMITATION NOTICES Qualified requesters may obtain copies of this report from DDC.			
11. SUPPLEMENTARY NOTES		12. SPONSORING MILITARY ACTIVITY Office of Naval Research	
13. ABSTRACT  A brief summary of recent major observations of magnetospheric boundary phenomena is provided with emphasis on Explorers 12 and 14 and IMP 1 measurements of charged particles and of magnetic fields in the sunward magnetopause region and within the tail of the magnetosphere in the anti-solar direction. Major observations of charged particles in these regions are summarized in several equatorial spatial distribution diagrams.			

14. KEY WORDS	LINK A		LINK B		LINK C	
	ROLE	WT	ROLE	WT	ROLE	WT
Observations of Magnetospheric Boundary Phenomena						

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